

Charge Injection and Transport in Conjugated Polymers

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Cornell University

Outline

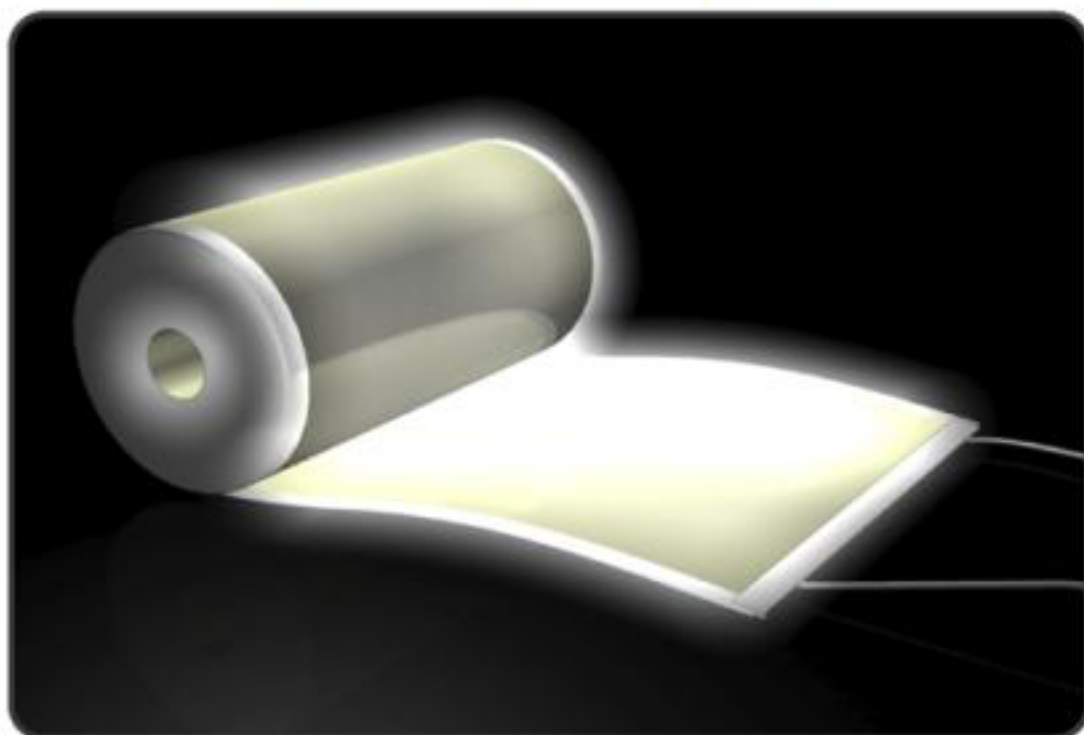
- Brief introduction to PLEDs
- Charge transport: theory and experiments
- Charge injection: theory and experiments
- Outlook



GE OLED Vision

“Lighting Wallpaper”

- Energy Efficient
- Low Cost
- Thin and Flexible



New design possibilities could change the way we think about lighting!

Early Progress

2001 - Illumination-Quality Light Possible



2002 - Scalable Large Area Architectures.

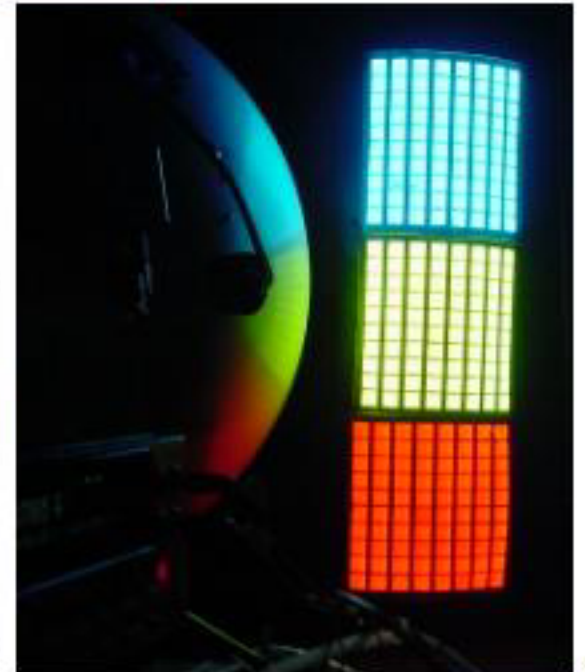
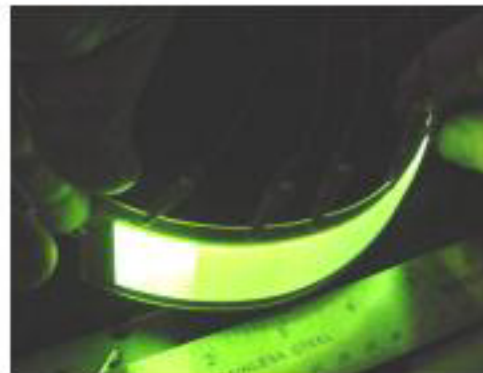
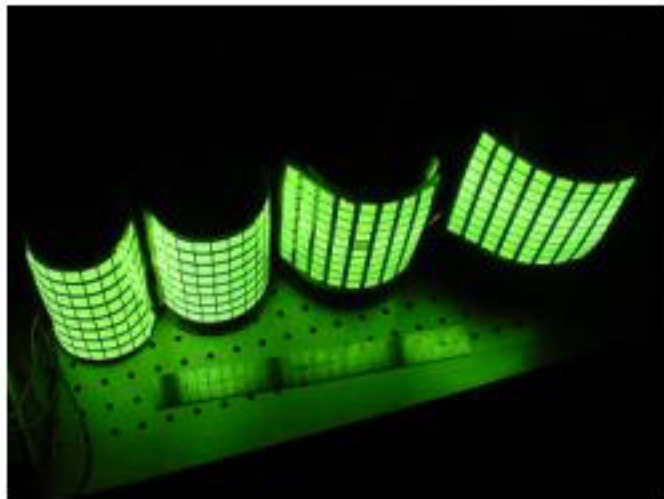
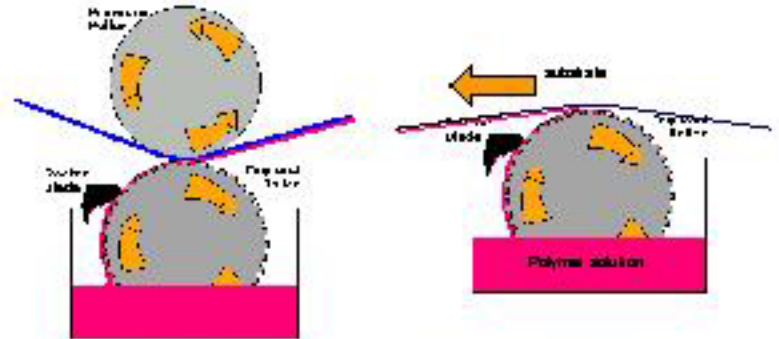
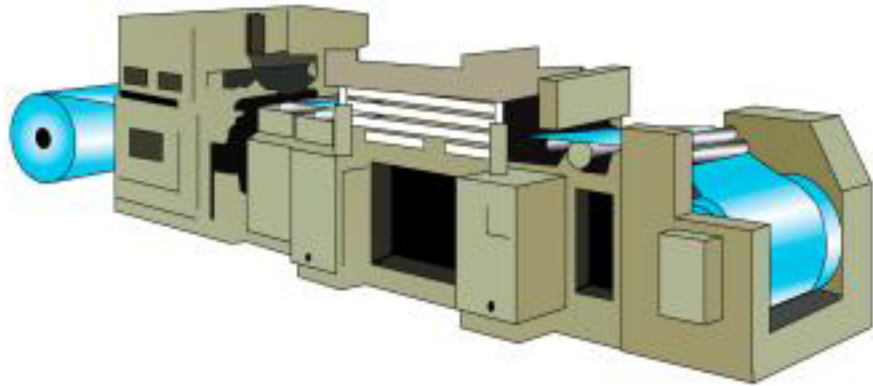


2003 - Incandescent Milestone

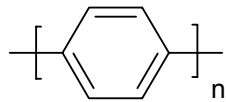


Current GE Focus

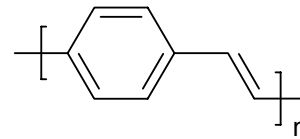
Developing Low Cost Manufacturing Infrastructure



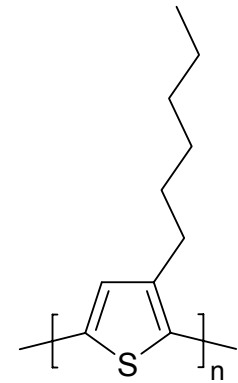
Chemical structure of conjugated polymers



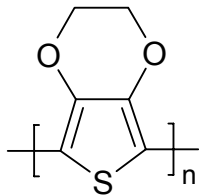
PPP



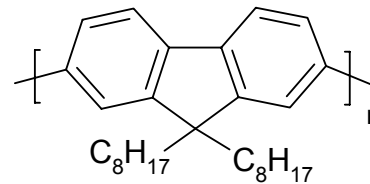
PPV



P3HT



PEDOT

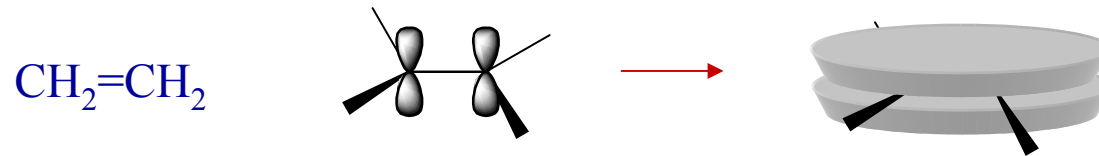


PFO



Carbon as a semiconductor

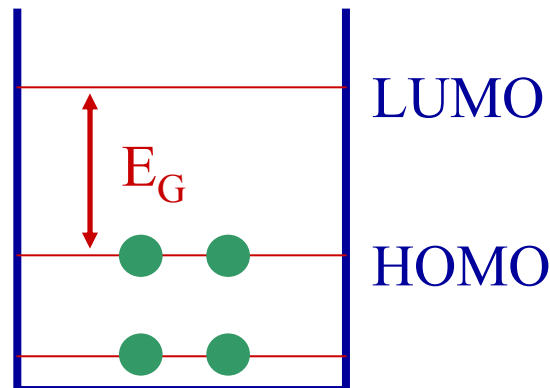
- Hybridization: sp^2 and p_z



- Particle in a box:

$$E_n = \frac{\hbar^2 \pi^2}{2mL^2} n^2$$

$$n=1,2,3,\dots$$



$$E_G \approx \frac{\hbar^2 \pi^2}{2maN}$$



Tuning of optical properties

Blue



Red



Table 1. Chemical structures and molecular weight characterization of regiospecific alkylated polythiophenes.

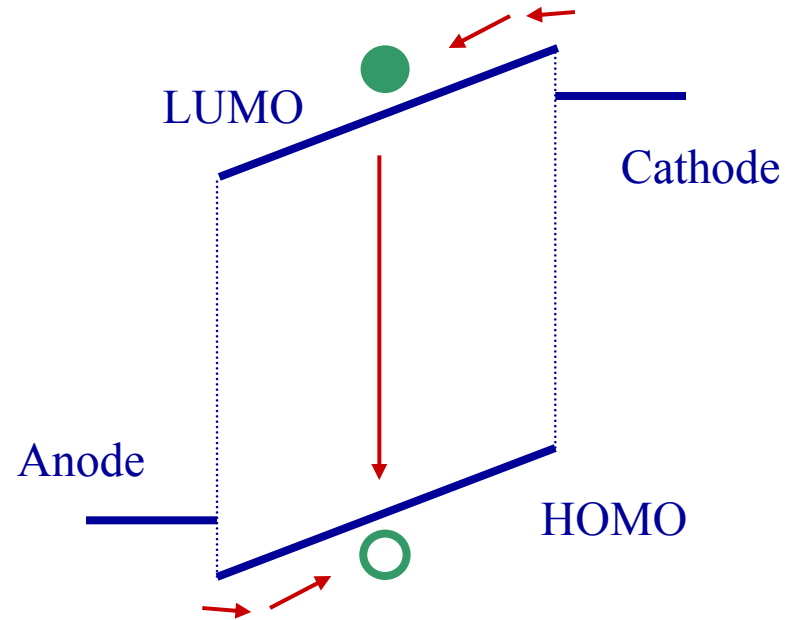
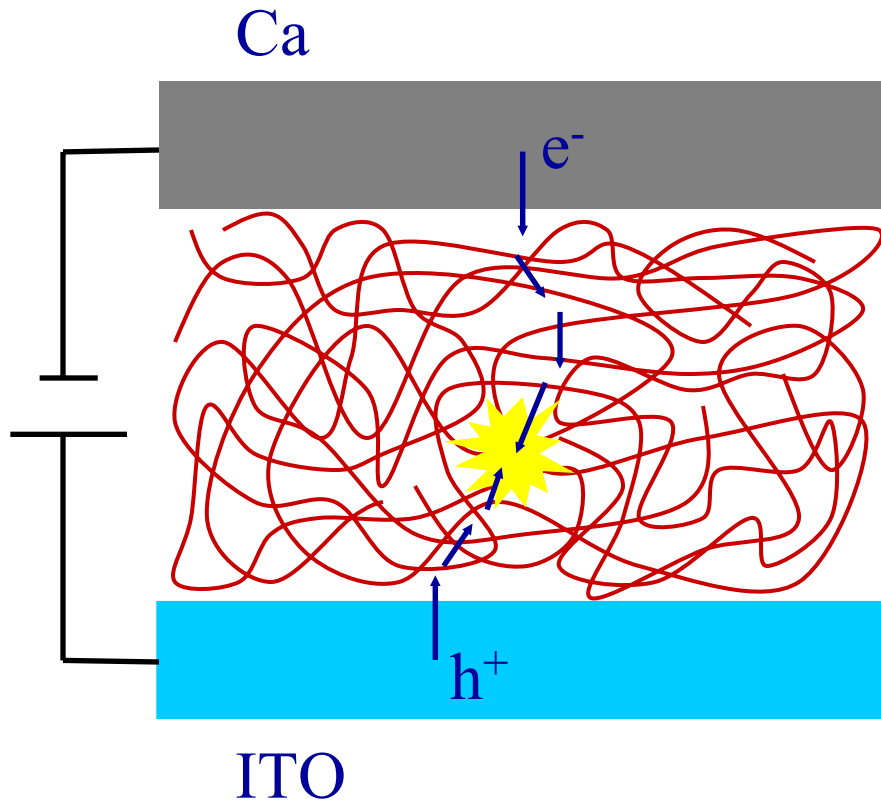
Polymer ^a	$M_n \times 10^{-4}$	M_w/M_n	
	I	1.7	2.3
	II	0.89	1.6
	III	4.2	2.7
	IV	8.5	2.1

[a] R is n-octyl. [b] Relative to polystyrene standards.



Covion

PLED structure and operation



Electroluminescence: charge injection,
charge transport, recombination



Outline

- Brief introduction to PLEDs
- Charge transport: theory and experiments
- Charge injection: theory and experiments
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Hierarchy of transport models

$$J = e \cdot n \cdot \mu \cdot E + e \cdot D \cdot dn/dx$$

First order correction:

Space charge effects

$$J = (9/8) \cdot \epsilon \cdot \epsilon_0 \cdot \mu \cdot V^2 / L^3$$

Disorder:

Energetics

Influence on mobility

Localized states

$$\mu = \mu(E, T)$$

Manifold filling:

Charge density dependence of mobility

$$\mu = \mu(n)$$

Charge generation:

Electric field dependence of charge density

$$n = n(E)$$



Space charge effects

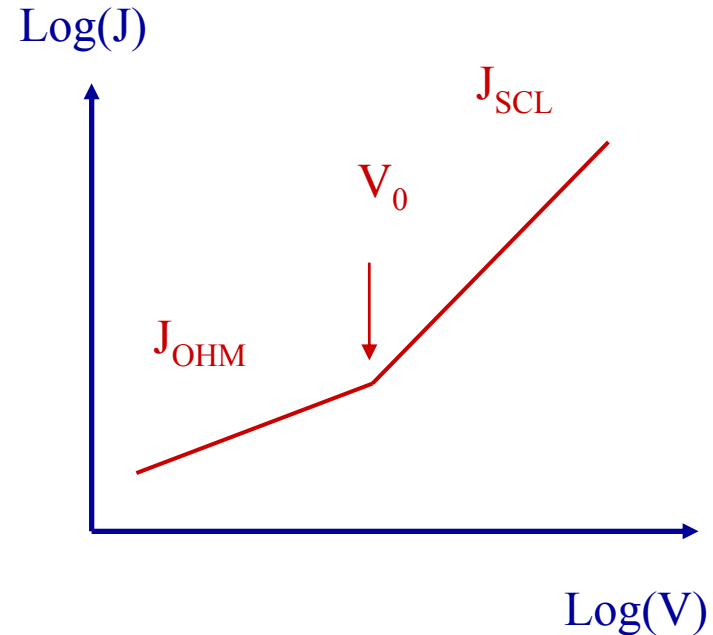
$$V_0 = (8/9) \cdot e \cdot N_0 \cdot L^2 / \epsilon \cdot \epsilon_0$$

Low voltages: Ohm's law

$$J_{\text{OHM}} = e \cdot N_0 \cdot \mu \cdot V / L$$

High voltages: Space charge limited current

$$J_{\text{SCL}} = (9/8) \cdot \epsilon \cdot \epsilon_0 \cdot \mu \cdot V^2 / L^3$$



Energetics of semiconductors

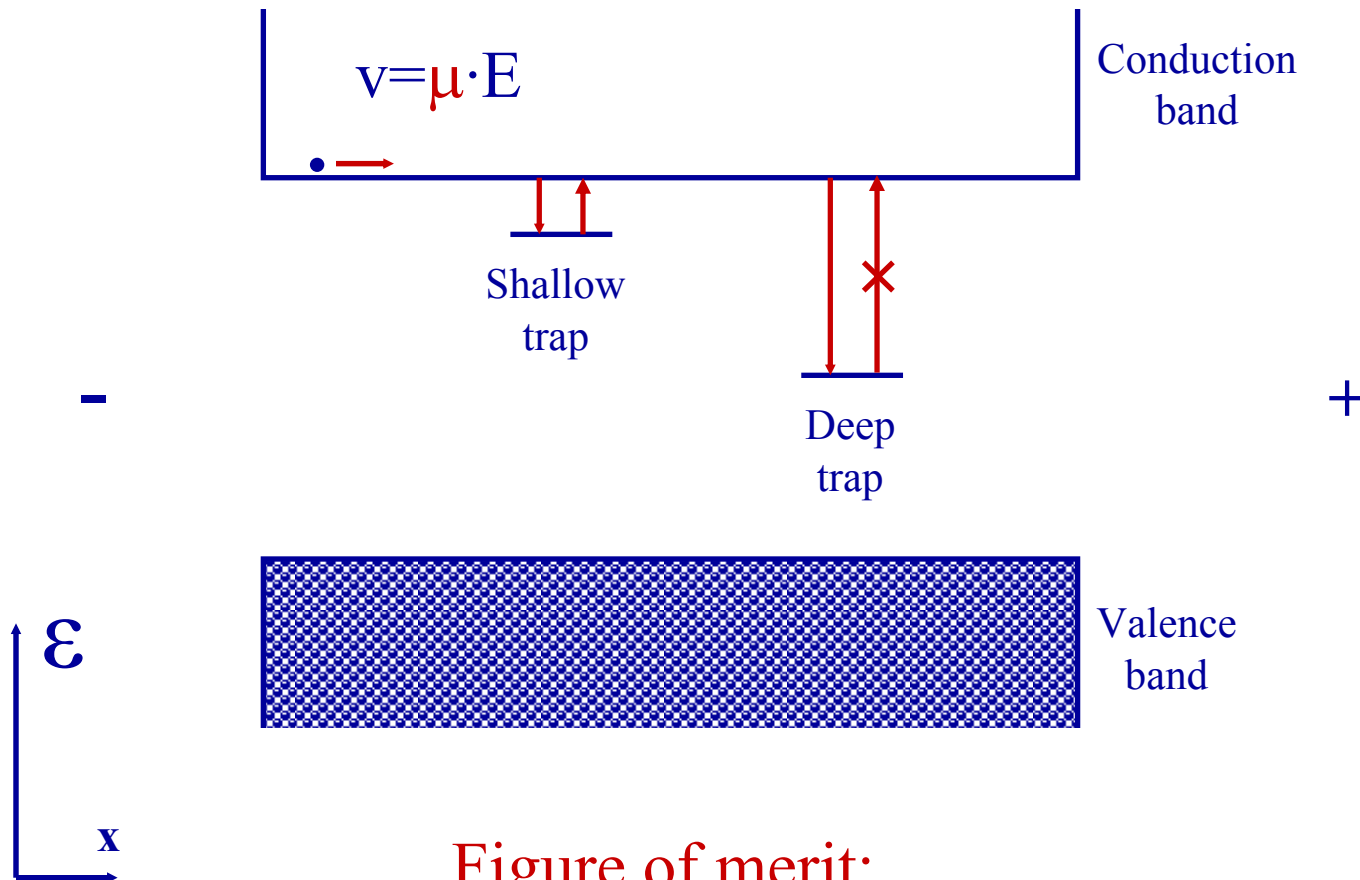
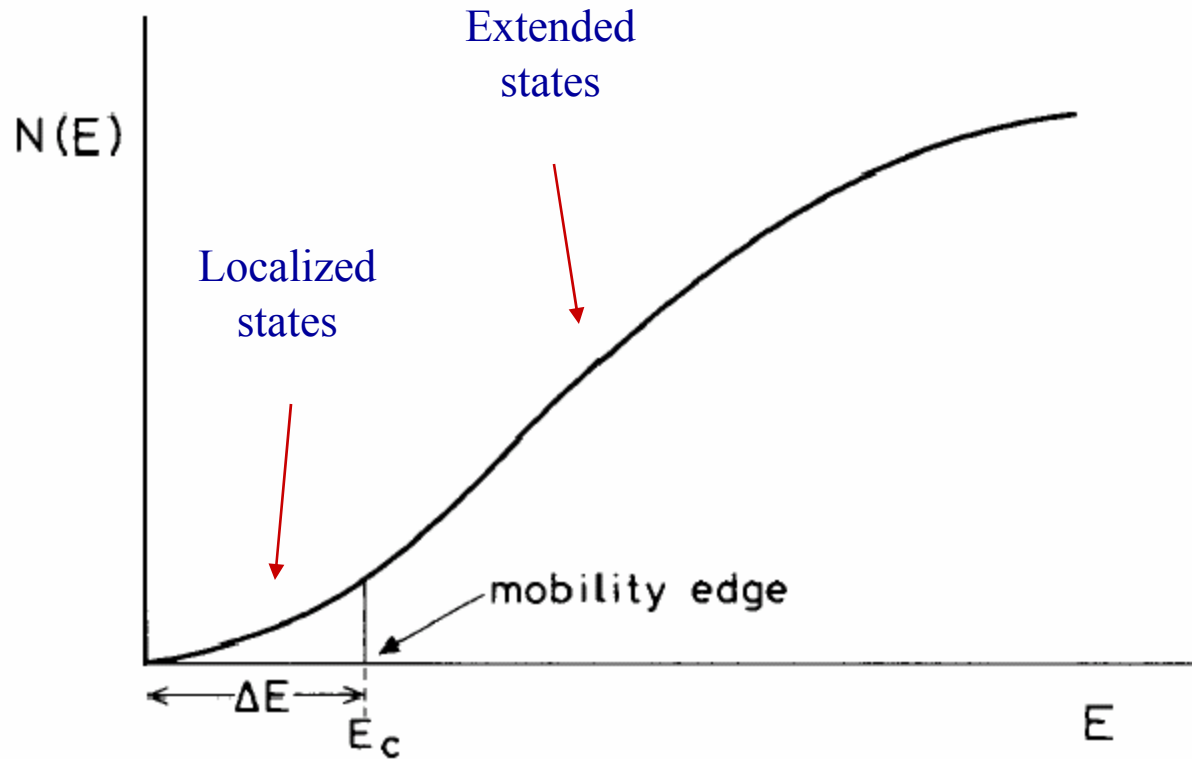


Figure of merit:
mobility, μ ($\text{cm}^2/\text{V}\cdot\text{sec}$)



Energetics of amorphous semiconductors

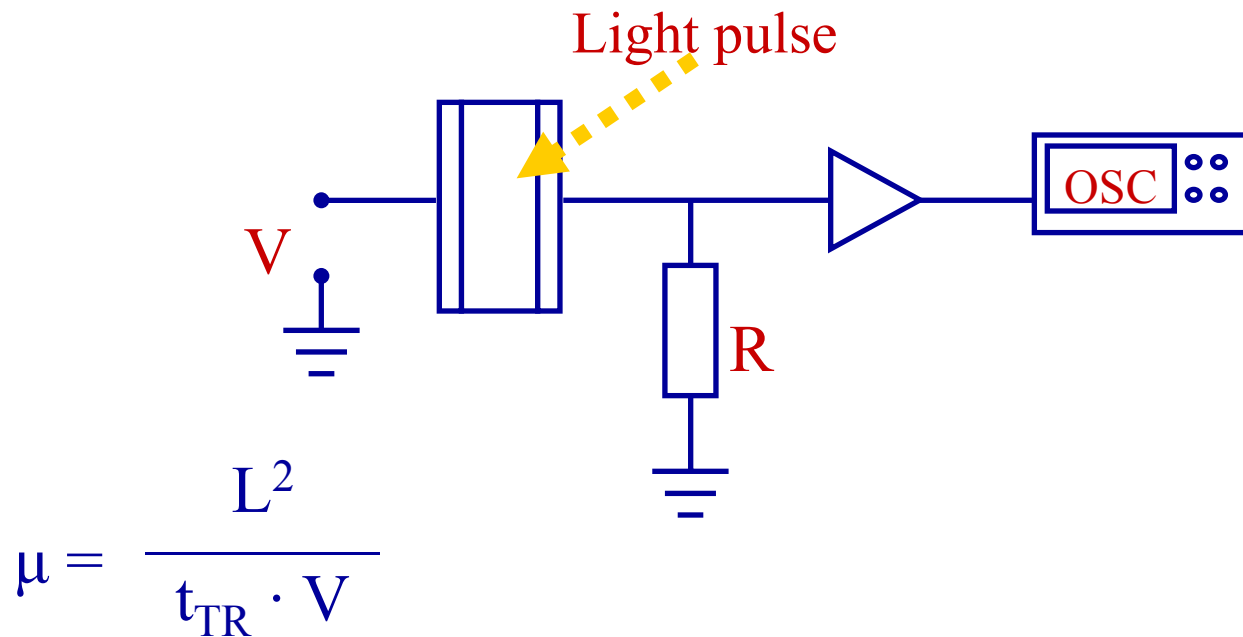


N. Mott, *Nobel Lecture* (1977)



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Time-of-flight (TOF)

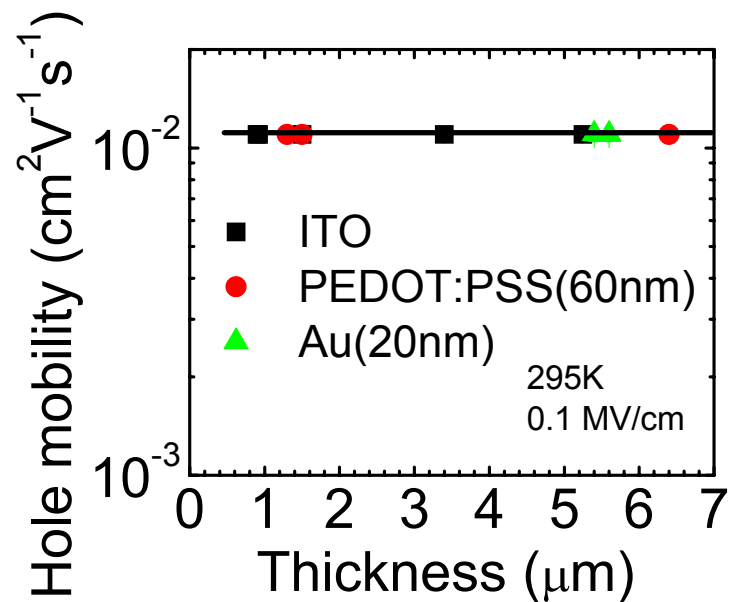
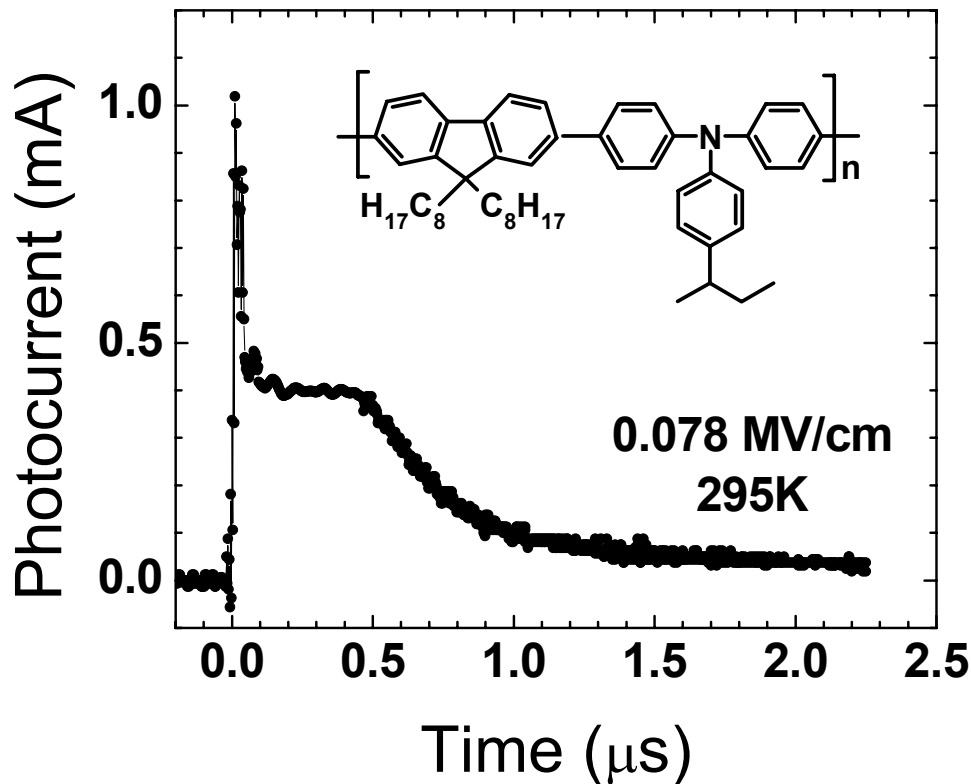


P. M. Borsenberger, D. S. Weiss, *Organic photoreceptors for Xerography* (Marcel Decker, Inc., New York, 1998).

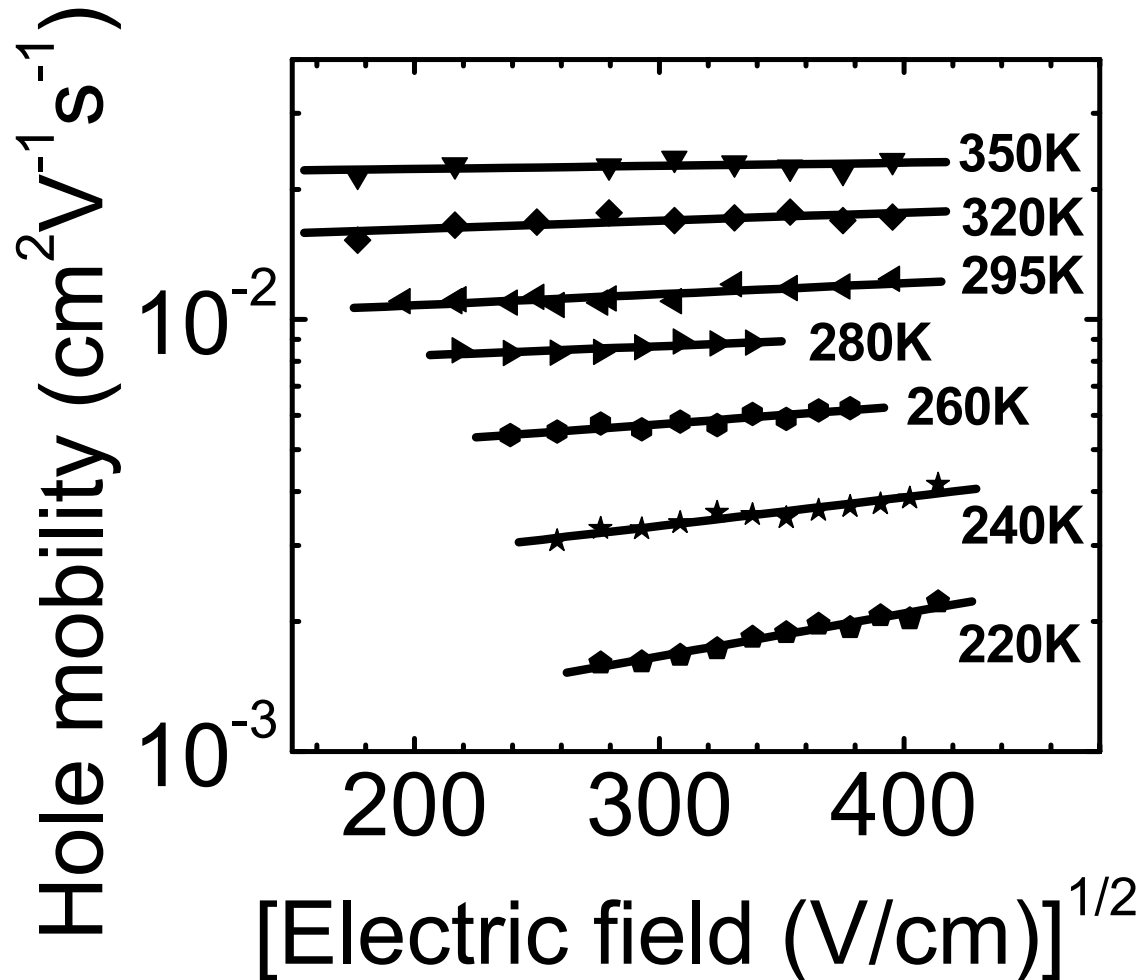


Non-dispersive hole transport in TFB

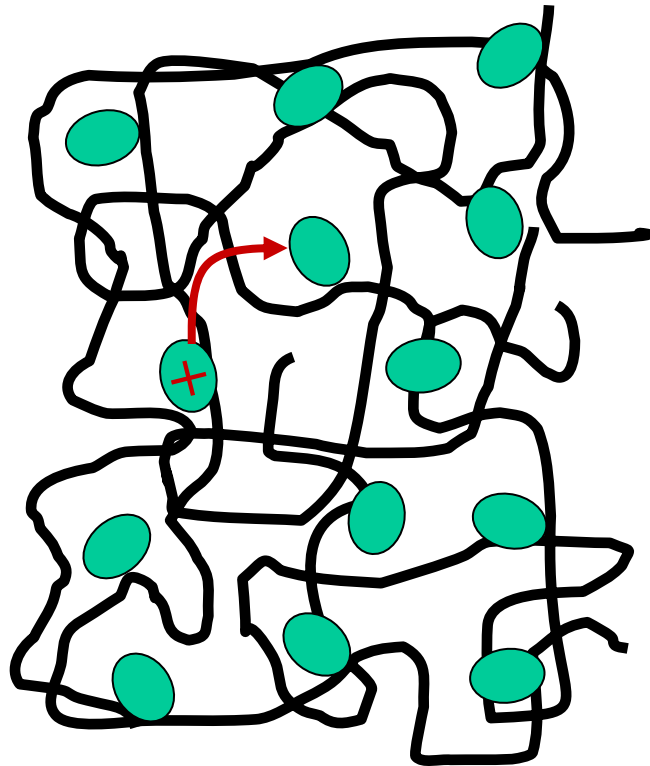
ITO/PEDOT:PSS(CH8000)/TFB(6.4 μm)/Al



Electric field dependence of mobility



Molecularly dispersed polymers



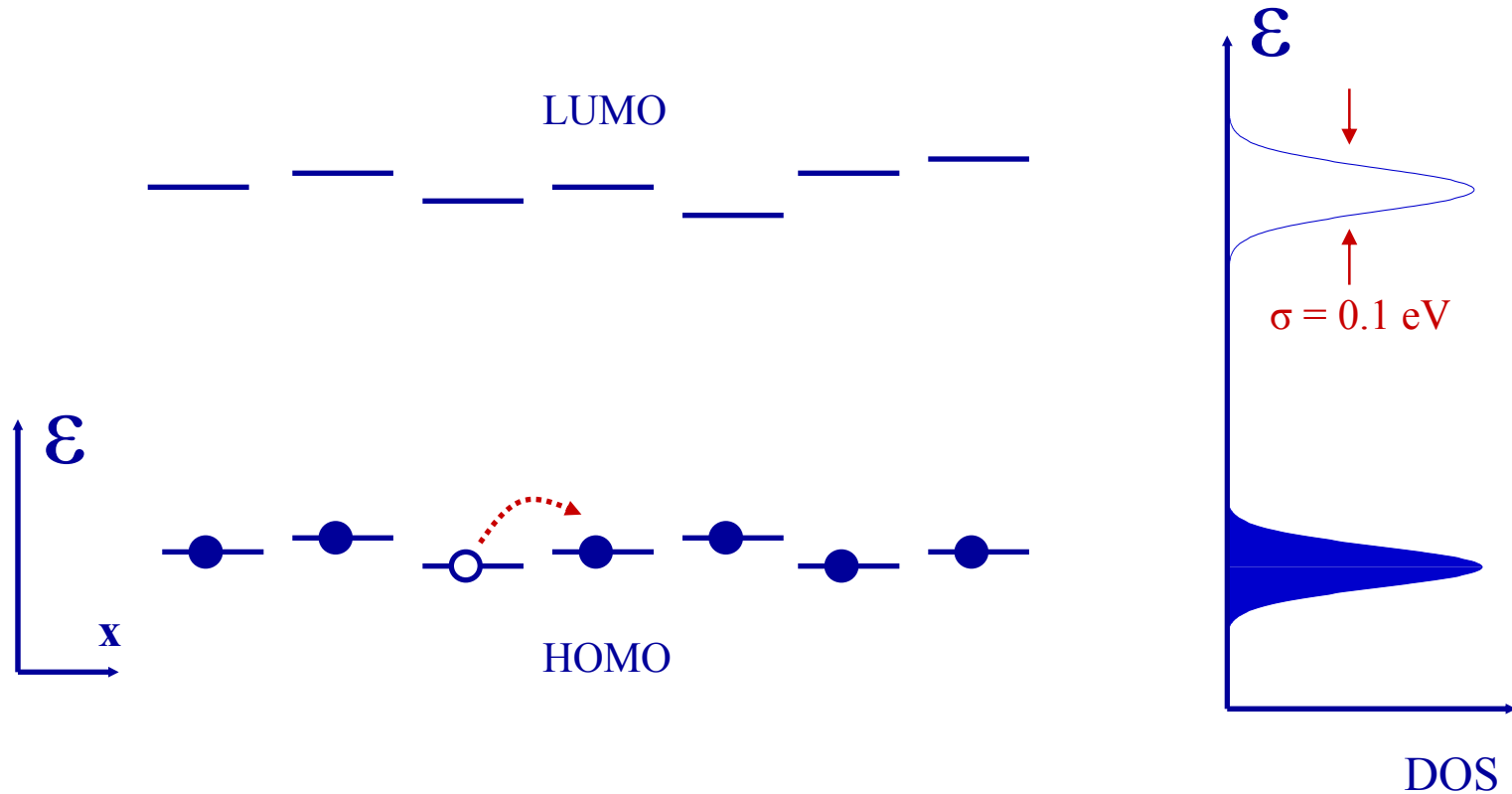
Solid solutions of
conjugated molecules
in inert host

Hopping sites are
well-defined

Control over the
average distance
between hopping sites



Gaussian disorder model



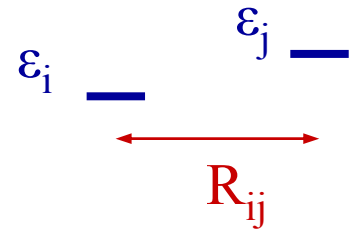
- Energetic disorder
- Positional disorder



Gaussian disorder model (II)

Density of states:

$$\text{DOS}(\varepsilon) = (2 \cdot \pi \cdot \sigma^2)^{-0.5} \cdot \exp[-(\varepsilon^2/2\sigma^2)]$$



Hopping rate:

$$v_{ij} = v_0 \cdot \exp-(2 \cdot \gamma \cdot a \cdot \Delta R_{ij}/R_{ij}) \cdot \begin{cases} \exp[-(\varepsilon_j - \varepsilon_i)/kT] ; \varepsilon_j > \varepsilon_i \\ 1 ; \varepsilon_j < \varepsilon_i \end{cases}$$

Mobility:

$$\mu = \mu_0 \cdot \exp[-(2\sigma/3kT)^2] \cdot \exp\{C \cdot [(\sigma/kT)^2 - \Sigma^2] \cdot E^{0.5}\}$$

Gaussian disorder model (III)

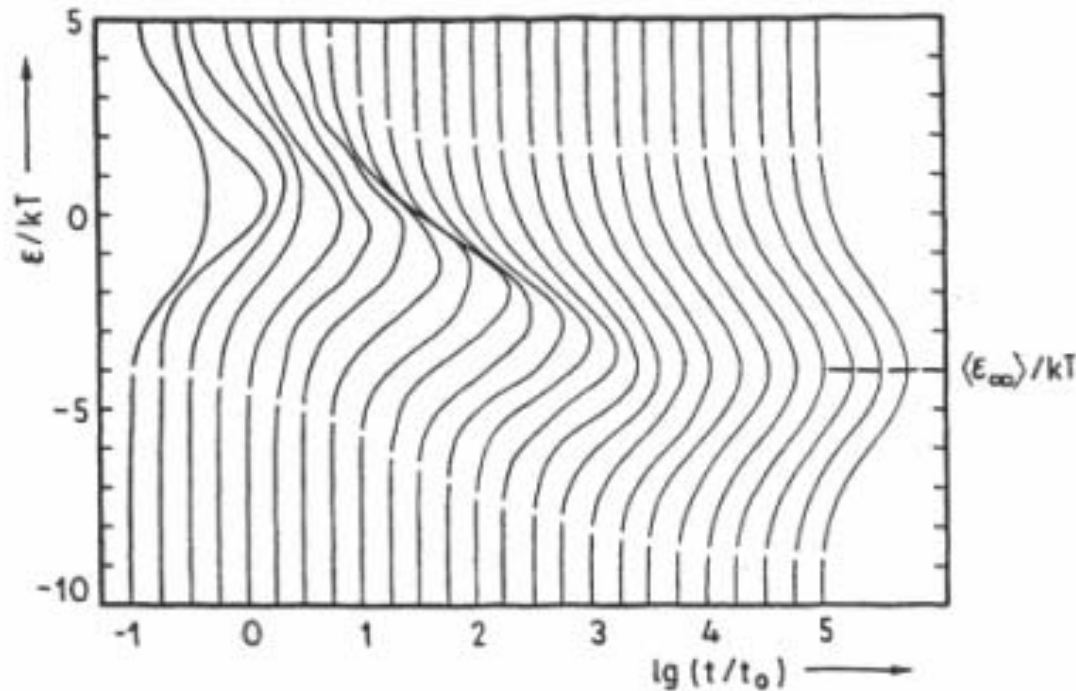


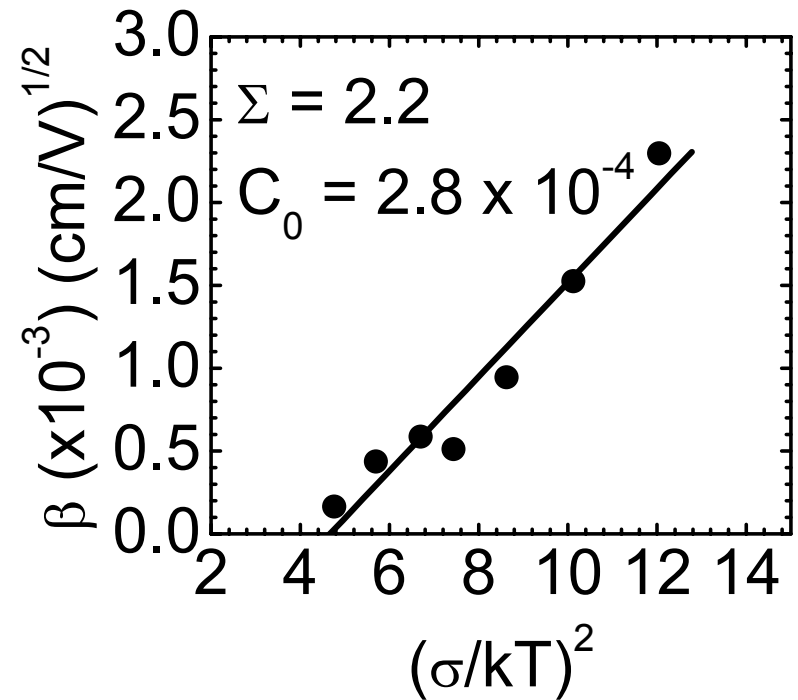
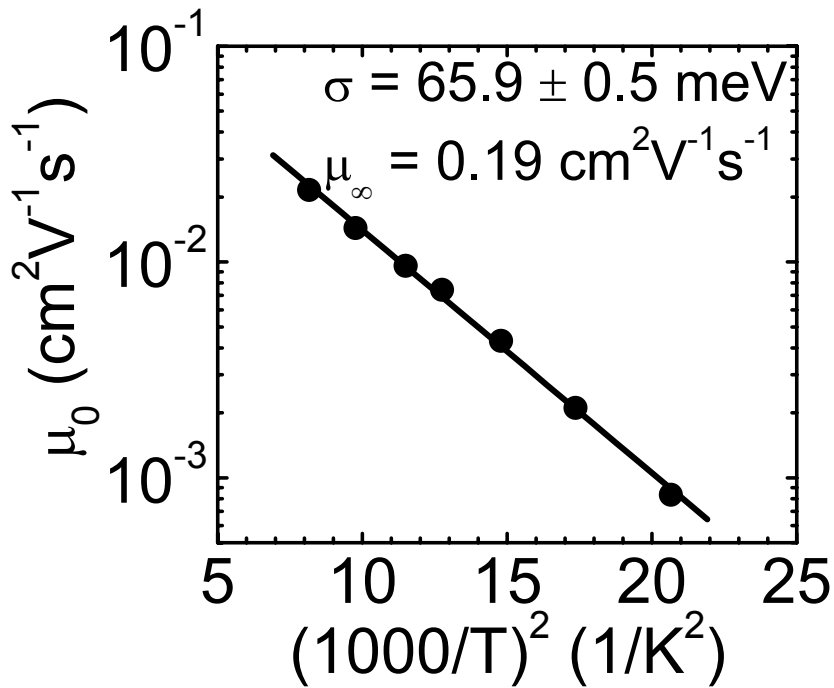
Fig. 2. Temporal evolution of the distribution of carrier energies in a Gaussian DOS of width $\hat{\sigma} = 2$. All profiles are broken at the same carrier density illustrating the different relaxation patterns for mobile and immobile carriers. ϵ_{∞} denotes the theoretical mean energy in the long-time limit

Carriers relax at:

$$\sigma^2/kT$$

$$\mu \sim \exp[-(\sigma/kT)^2]$$

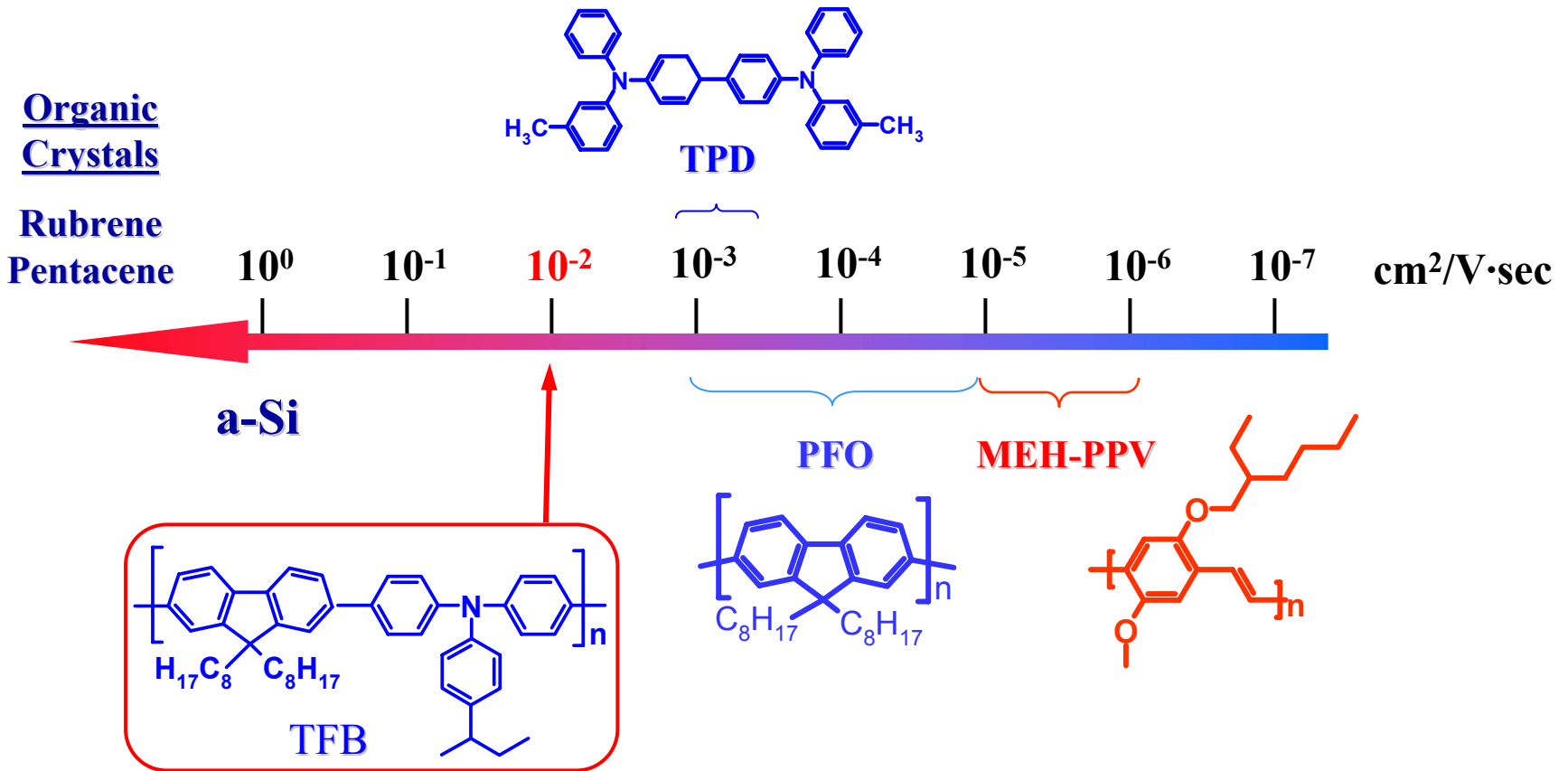
Electric field dependence of mobility (II)



$$\mu = \mu_0 \cdot \exp[-(2\sigma/3kT)^2] \cdot \exp\{C \cdot [(\sigma/kT)^2 - \Sigma^2] \cdot E^{0.5}\}$$



High mobility in conjugated polymers



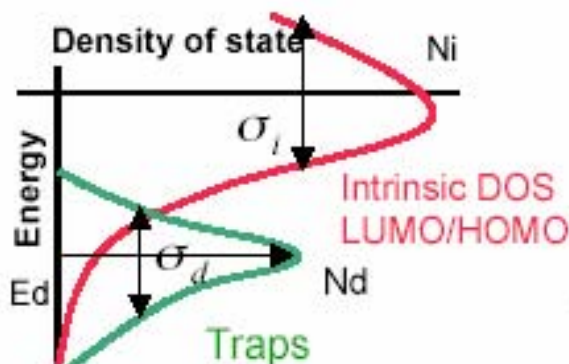


Physical Models for Analysis of Electrical Characteristics for Organic Devices

- Hopping Model

- Effective Transport Energy

- The effective carrier transport energy (E_{tr}) is calculated from^{1,2}:



$$\int_{-\infty}^{E_{tr}} g(E)(E_{tr} - E)^3 dE = \frac{6\beta}{\pi} (\gamma kT)^3$$

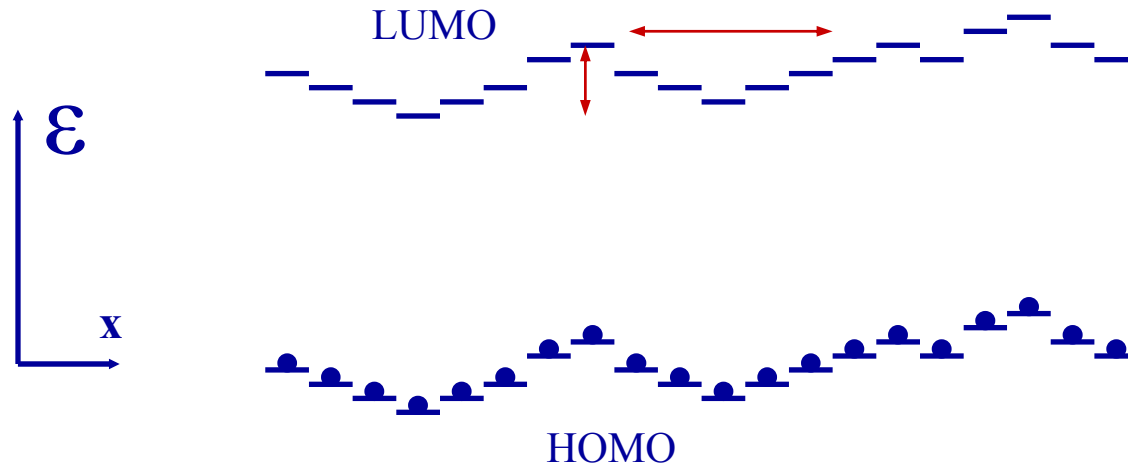
$$g(E) = \frac{N_i}{\sqrt{2\pi}\sigma_i} \exp\left(\frac{-E^2}{2\sigma_i^2}\right) + \frac{N_d}{\sqrt{2\pi}\sigma_d} \exp\left(\frac{-(E + E_d)^2}{2\sigma_d^2}\right)$$

Where $g(E)$ is the DOS distribution, N_i is the total intrinsic state density, N_d is the total dopant state density, σ_i is the intrinsic Gaussian DOS width, σ_d is the dopant Gaussian DOS width, E_d is the energy shift, γ is 1/carrier localization radius, β is the percolation constant, E is the band energy, k is Boltzmann's constant and T is the lattice temperature

¹Charge carrier mobility in doped disordered organic semiconductors" - V.I. Arkhipov, P. Heremans, E.V. Emelianova, G.J. Adriaenssens, H. Bassler, Journal of Non-Crystalline Solids, 338-340, pp 603-603, 2004.

²Charge carrier mobility in doped semiconducting polymers" - V.I. Arkhipov, P. Heremans, E.V. Emelianova, G.J. Adriaenssens, H. Bassler, Applied Physics Letters, Vol. 82, No. 19, pp 3245-3247, 2003.

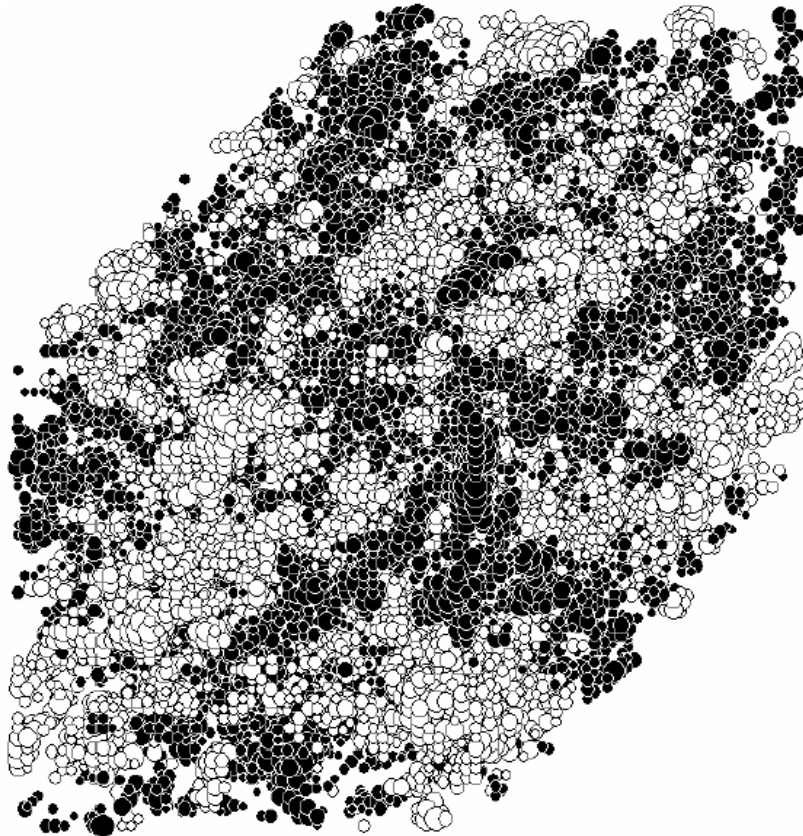
Correlated disorder model



Deeper valleys are also wider.

$$\mu = \mu_0 \cdot \exp[-(\sigma/kT)^2 + 2 \cdot (\sigma/kT) \cdot (e \cdot a \cdot E/kT)^{0.5}]$$

Correlated disorder model (II)



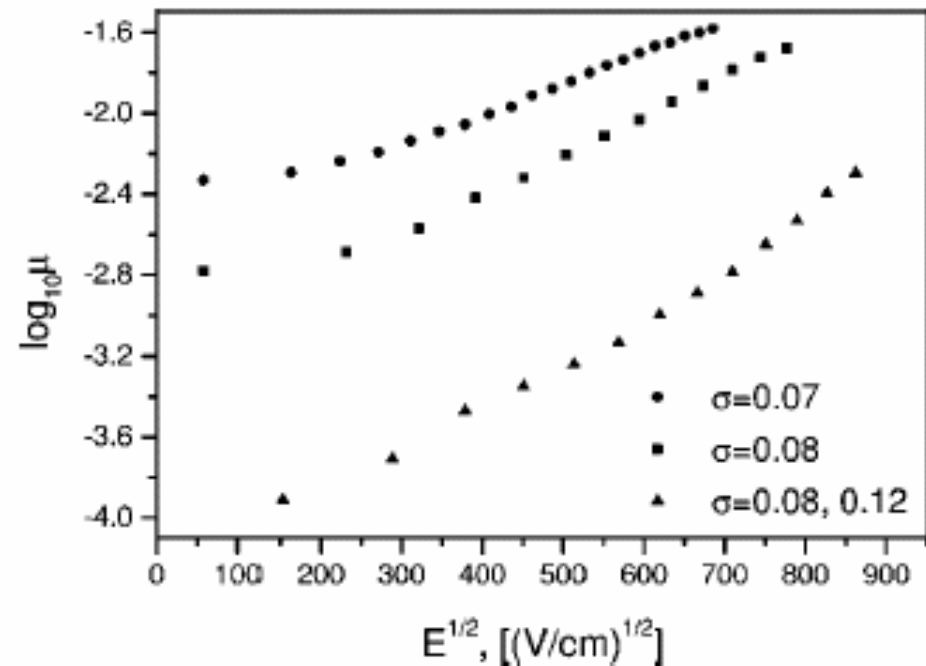
Black/white:
sites with energy
above/below the
mean.

Correlated disorder in polymers

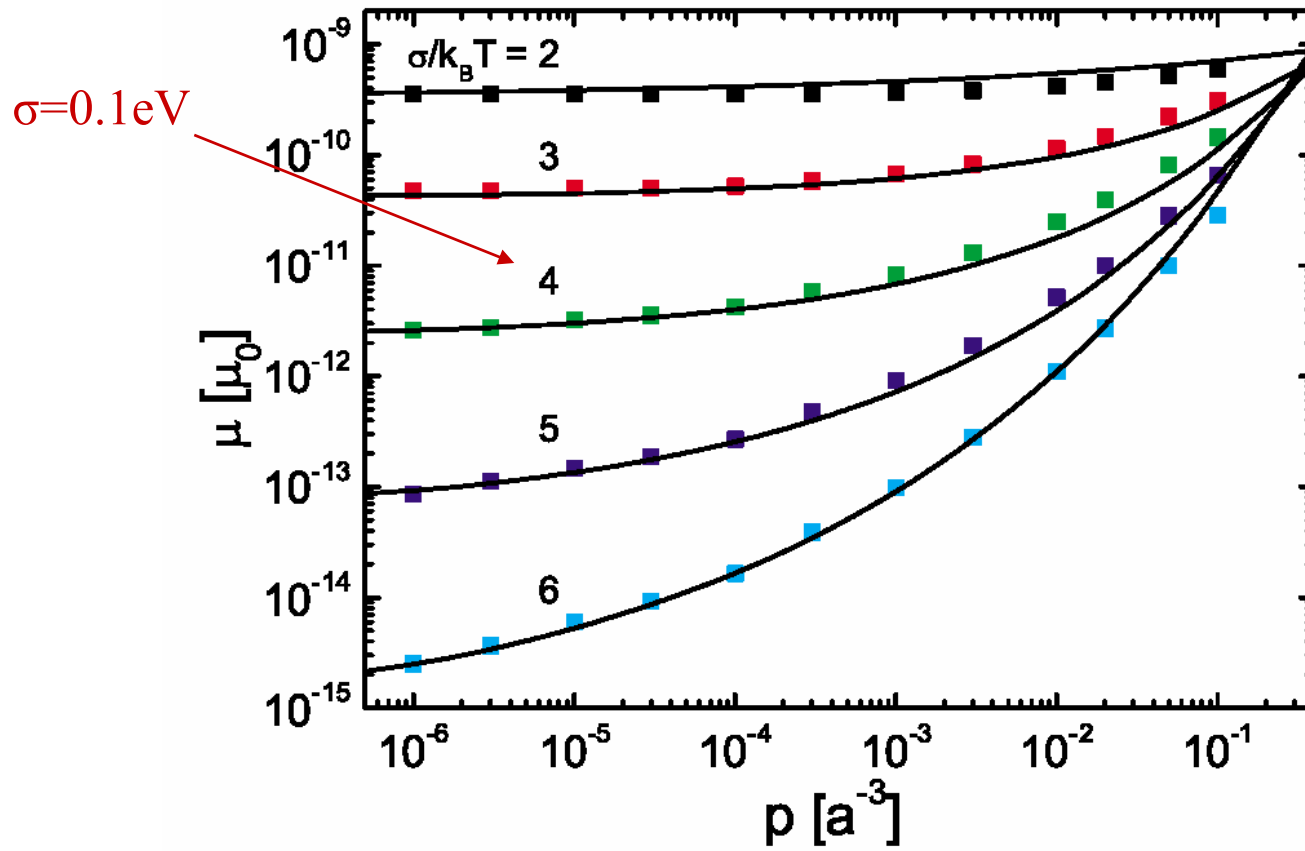
Correlations in site energy arise from fluctuations in:

-Conjugation length

-Density

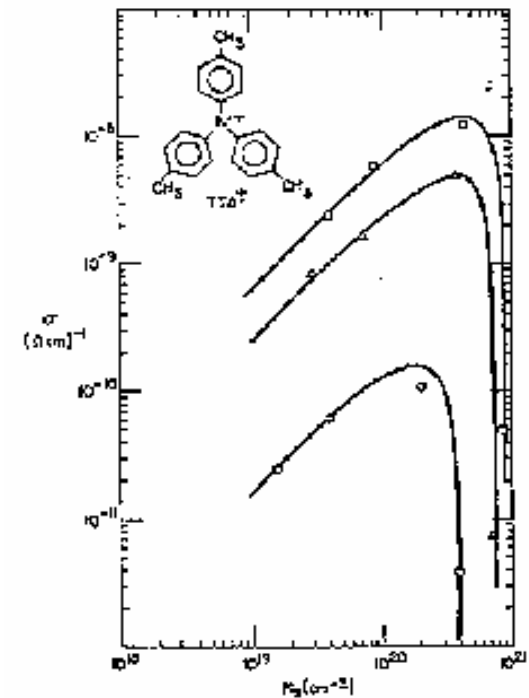
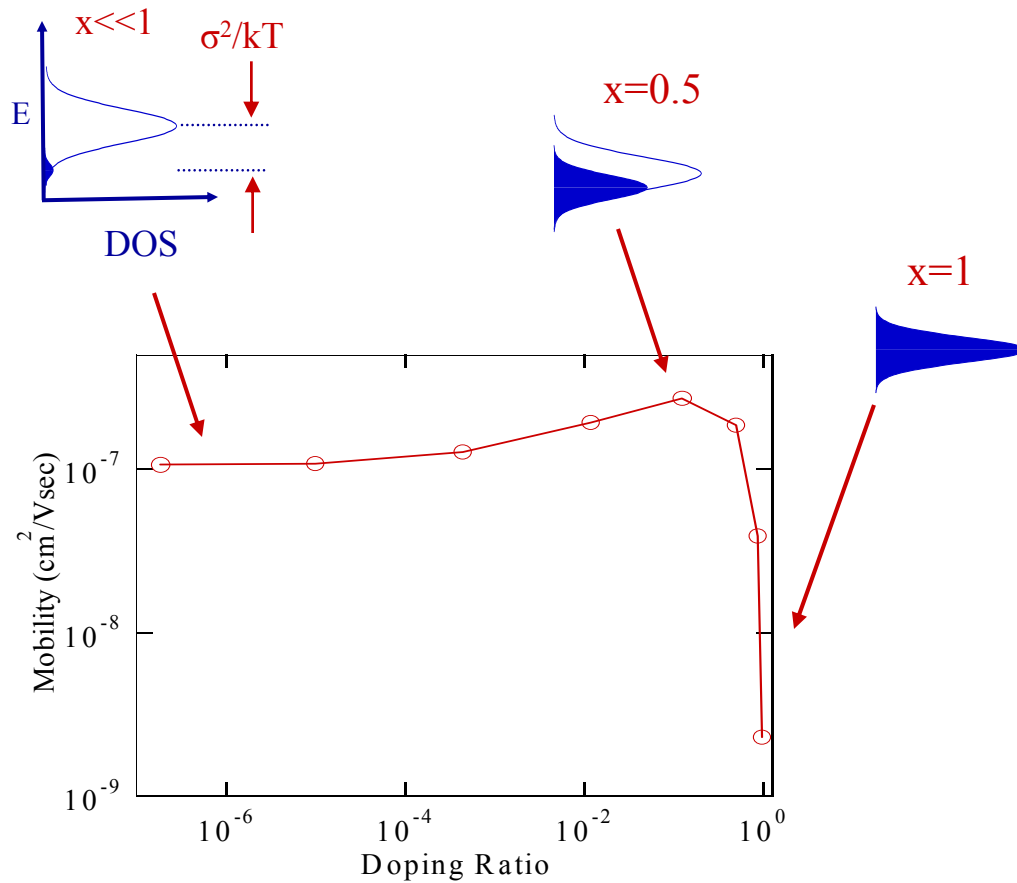


Mobility vs. charge density



W.F. Pasveer, J. Cottaar, C. Tanase, R. Coehoorn, P.A. Bobbert, P.W.M. Blom, D.M. de Leeuw, and M.A.J. Michels, *Phys. Rev. Lett.* **94**, 206601 (2005).

Manifold filling



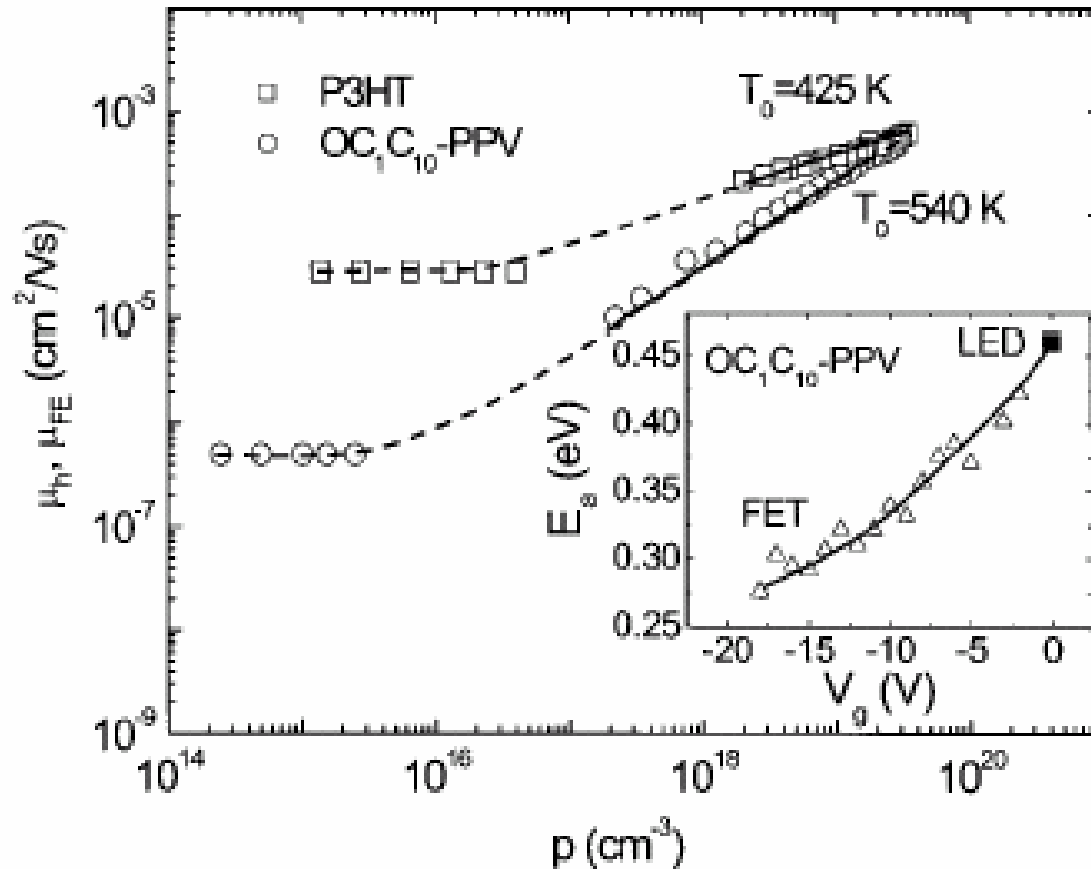
A. Troup *et al.*, *J. Non-Crystalline Solids* **35**, 151 (1980).

Y. Shen, K. Diest, M.H. Wong, B.R. Hsieh, D.H. Dunlap, and G.G. Malliaras, *Phys. Rev. B* **68**, 81204(R) (2003).



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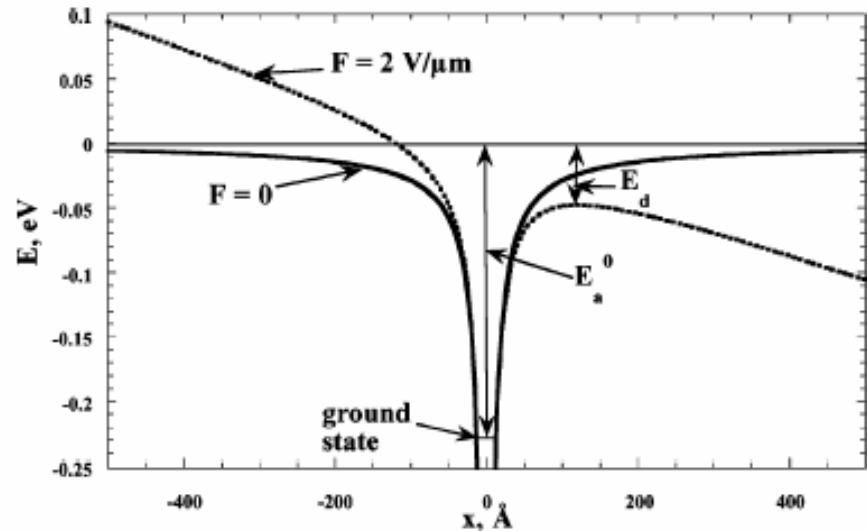
Mobility vs. charge density (II)



Charge density vs. electric field

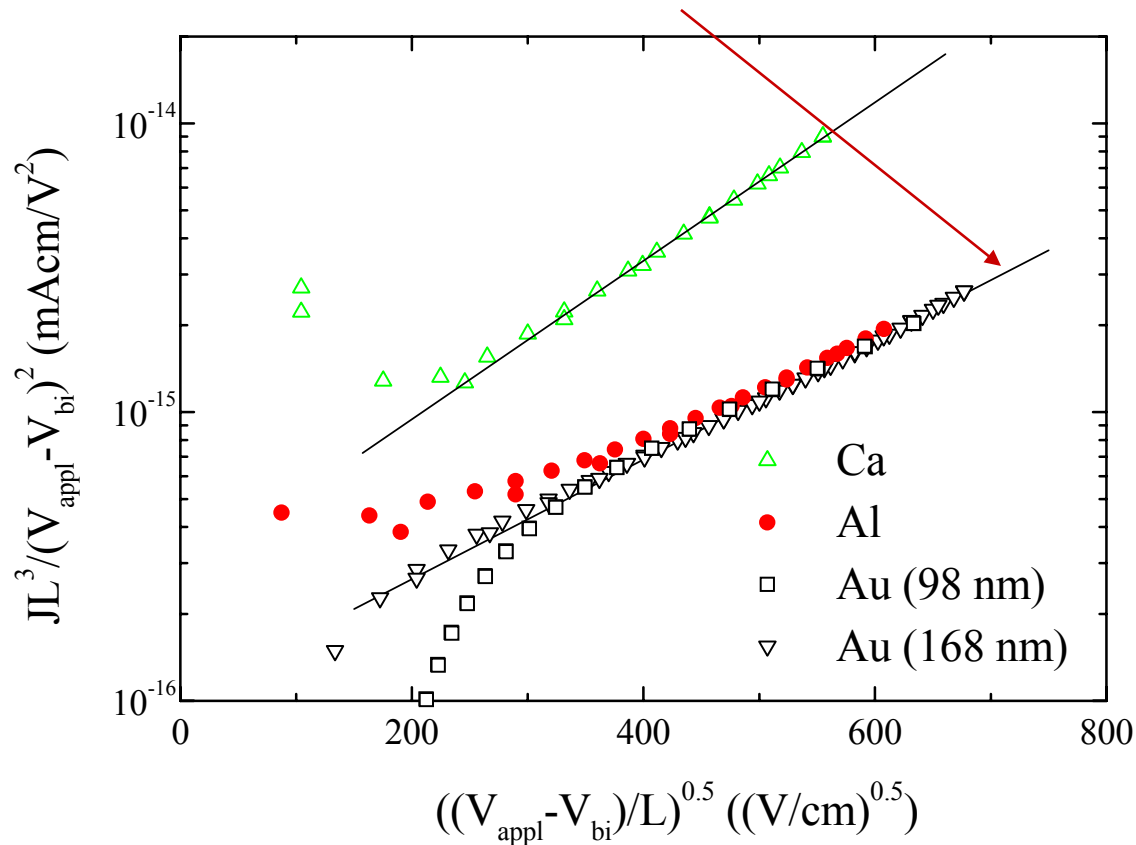
Field ionization of impurities leads to $n=n(E)$

$$J = q \mu_n F n_d \exp\left(\frac{\left(-E_a^0 + \left(\frac{q^3}{\pi \epsilon_0 \epsilon}\right)^{1/2} F^{1/2}\right)}{kT}\right)$$



Charge density vs. electric field (II)

$$J_{\text{SCL}} \approx (9/8)\epsilon\epsilon_0\mu_0 V^2 \exp[0.89(V/E_0 L)^{0.5}]/L^3$$



Charge transport

First order correction:

Space charge effects

$$J = (9/8) \cdot \epsilon \cdot \epsilon_0 \cdot \mu \cdot V^2 / L^3 \quad \checkmark$$

Disorder:

Energetics

Localized states ✓

Influence on mobility

$$\mu = \mu(E, T) \quad \checkmark$$

Manifold filling:

Charge density dependence of mobility

$$\mu = \mu(n) \quad \checkmark$$

Charge generation:

Electric field dependence of charge density

$$n = n(E) \quad ?$$



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- Charge injection: theory and experiments
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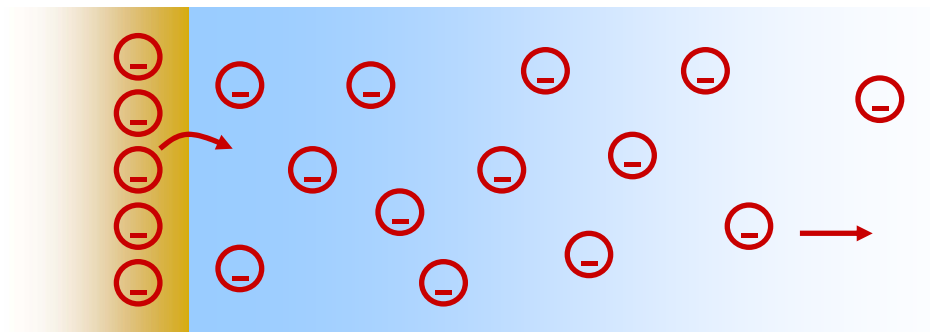


Injection vs. transport



Water hose
and valve

Is the flow limited by the valve or the hose?

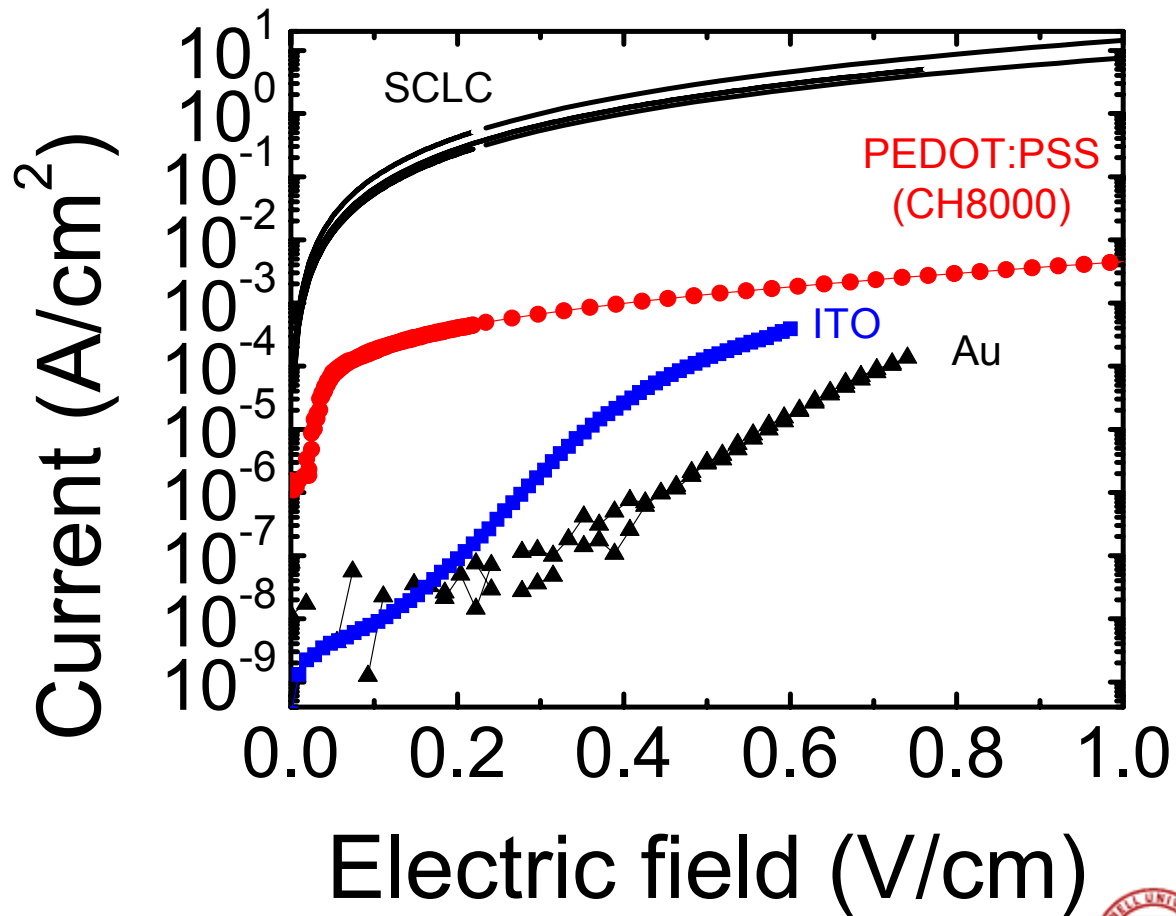


Semiconductor
contacts

Is the current limited by injection or transport?



Hole injection in TFB



Hierarchy of injection models

Mechanism:

Thermionic emission

$$J = A \cdot \exp(-\phi/kT)$$

Tunneling

$$J = A \cdot E^2 \cdot \exp(-B \cdot \phi^{3/2}/E)$$

First order corrections:

Barrier lowering

$$J \sim \exp(E^{0.5})$$

Recombination with image force

$$J \sim \mu$$

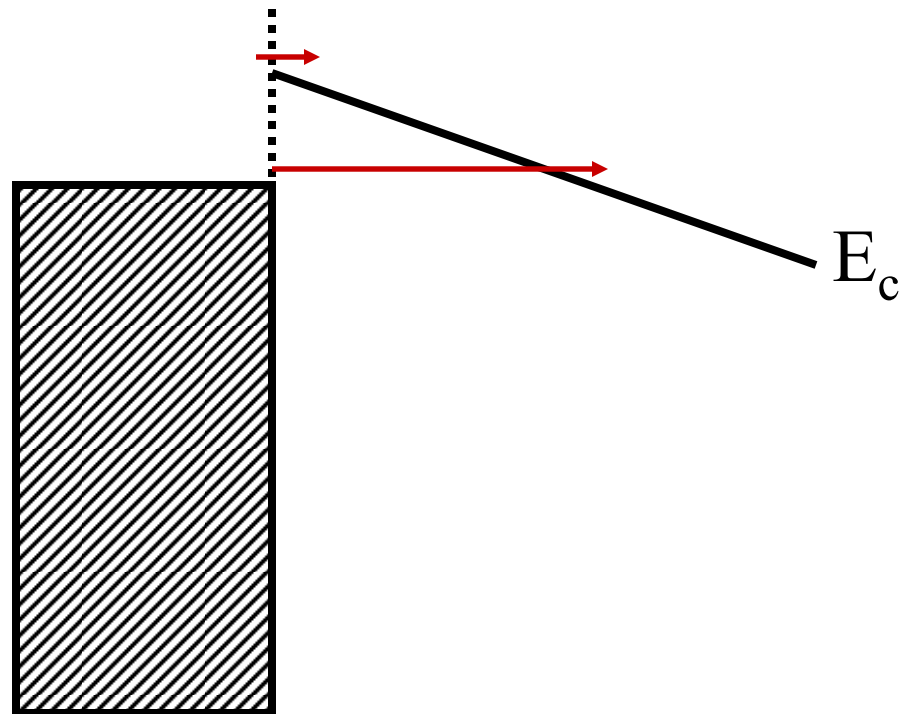
Disorder:

Gaussian disorder

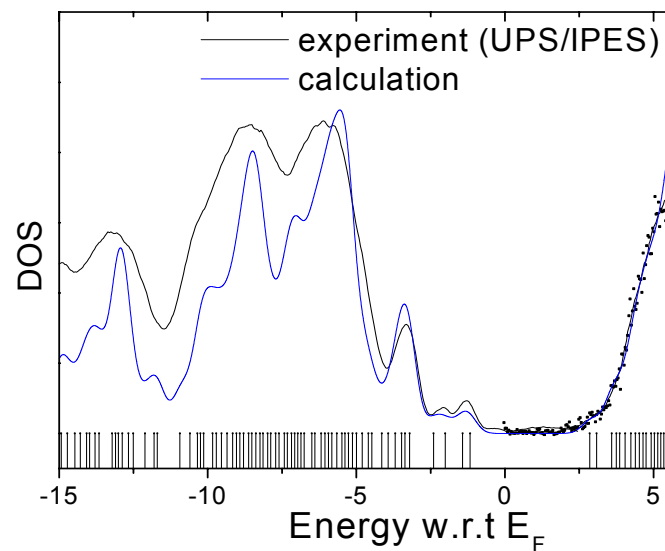
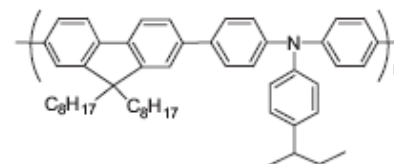
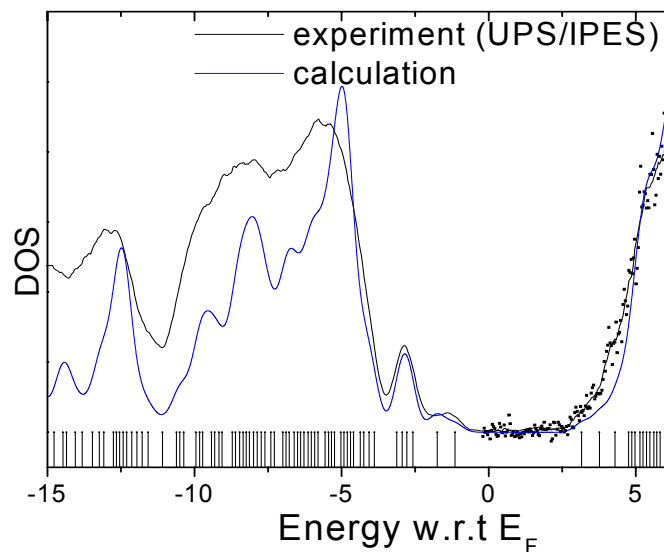
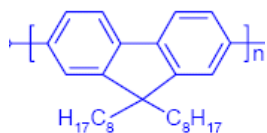
$$J \sim \exp(E^{0.5})$$



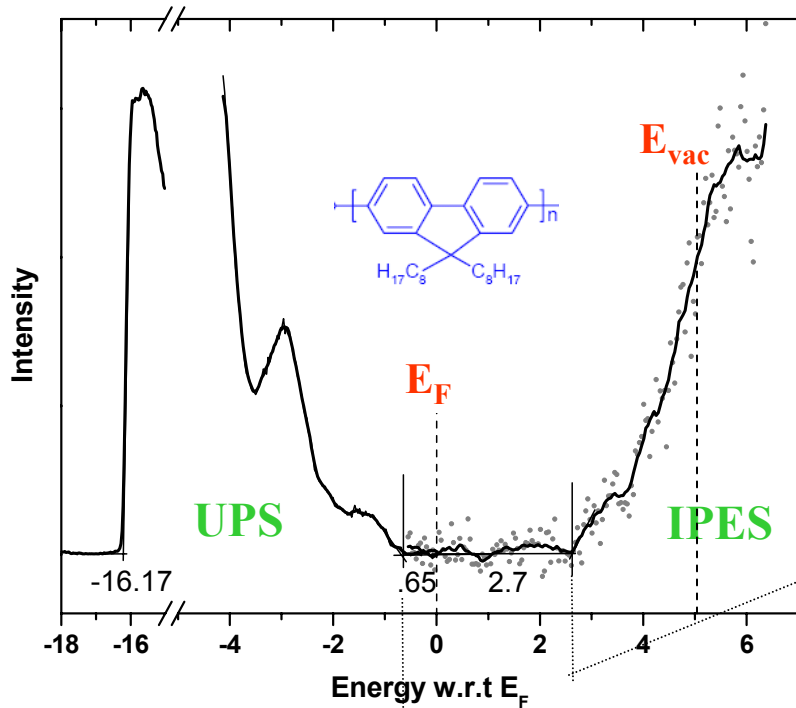
Thermionic emission and tunneling



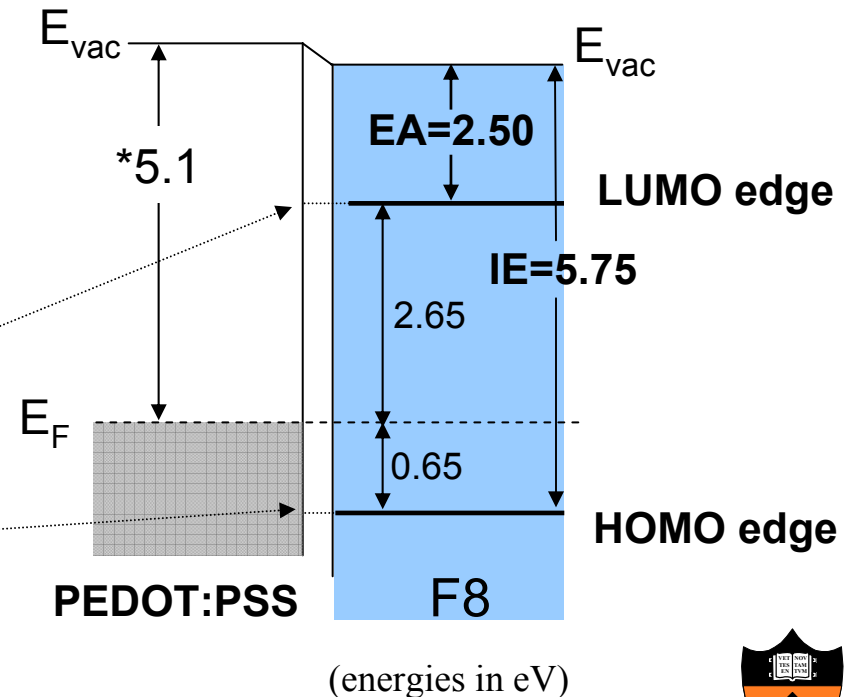
Energetics of conjugated polymers



Energetics at the contact



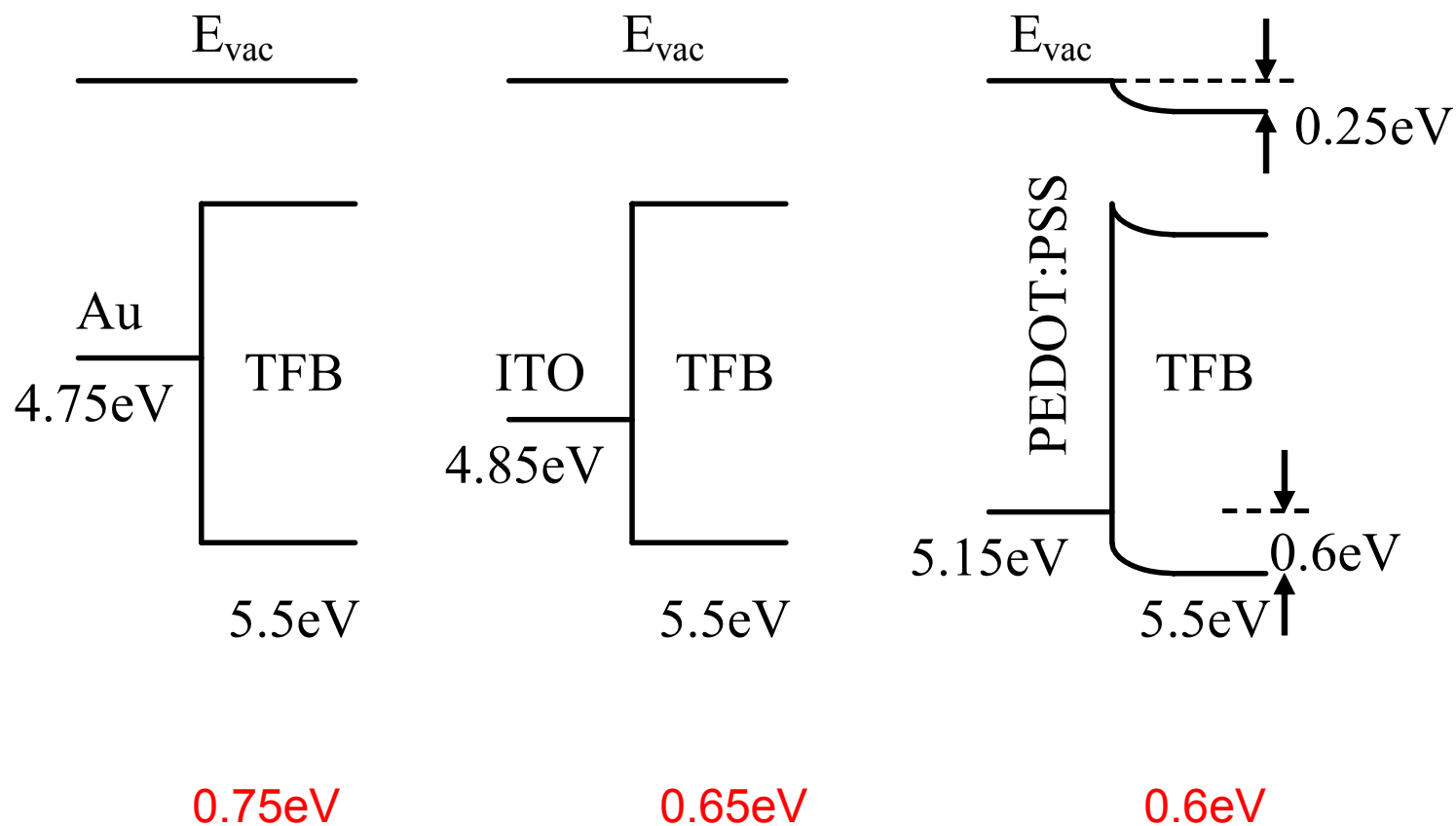
Ionization energy (IE) and electron affinity (EA) measured from HOMO and LUMO edges w.r.t vacuum level



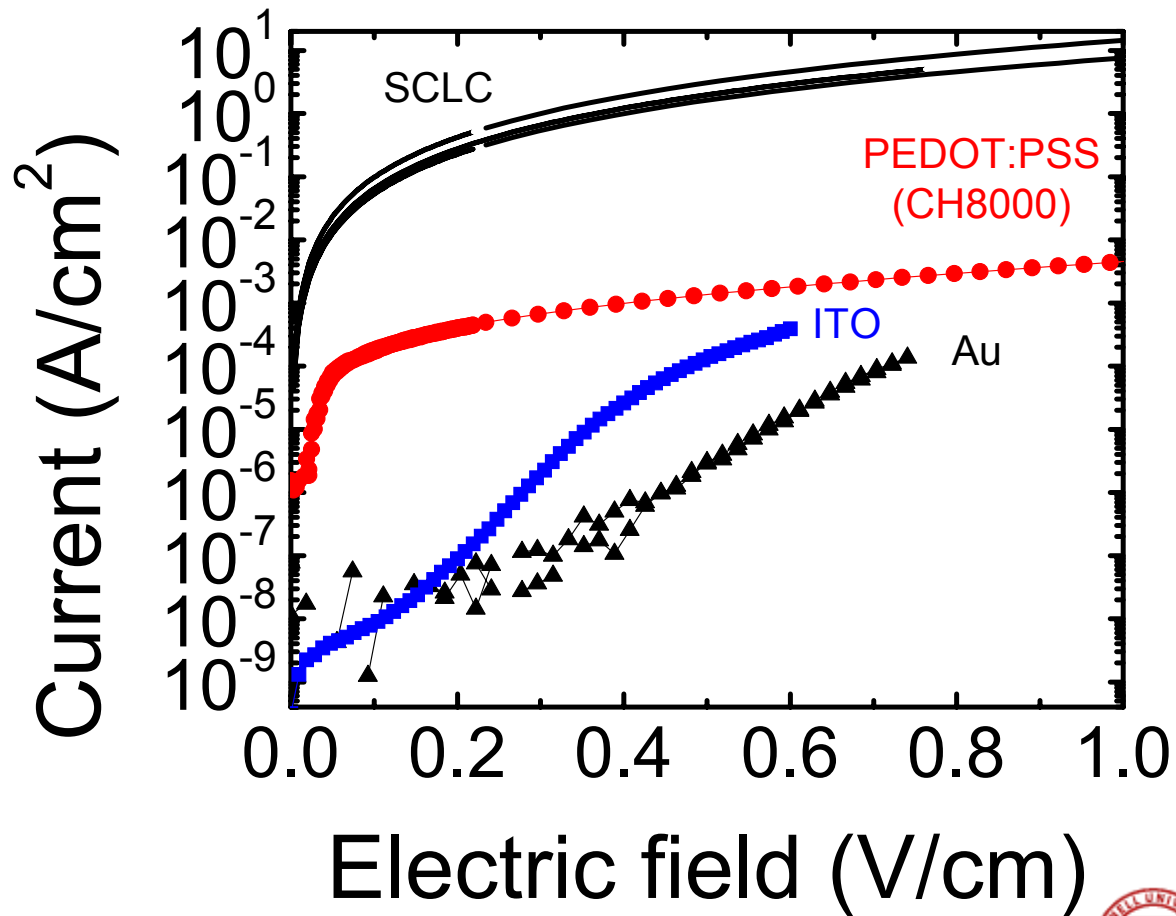
* Work function of PEDOT:PSS substrate measured on separate sample produced in same batch.



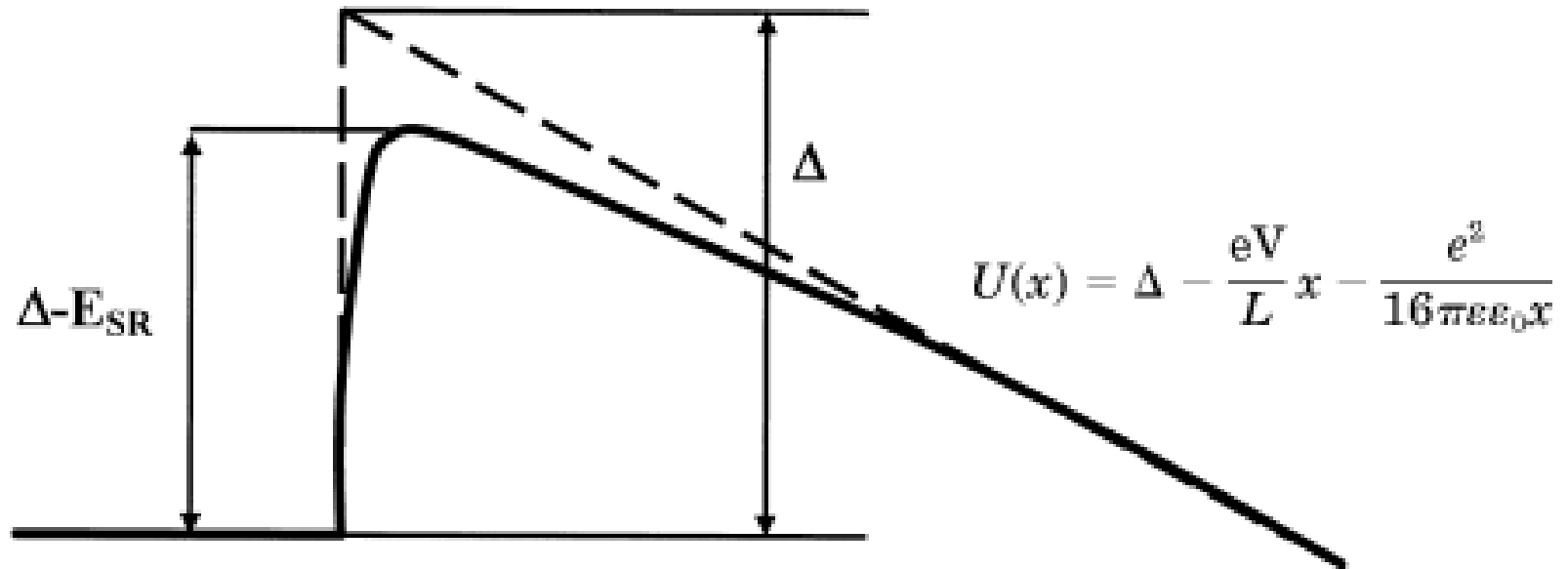
Hole injection barriers for TFB contacts



Hole injection in TFB



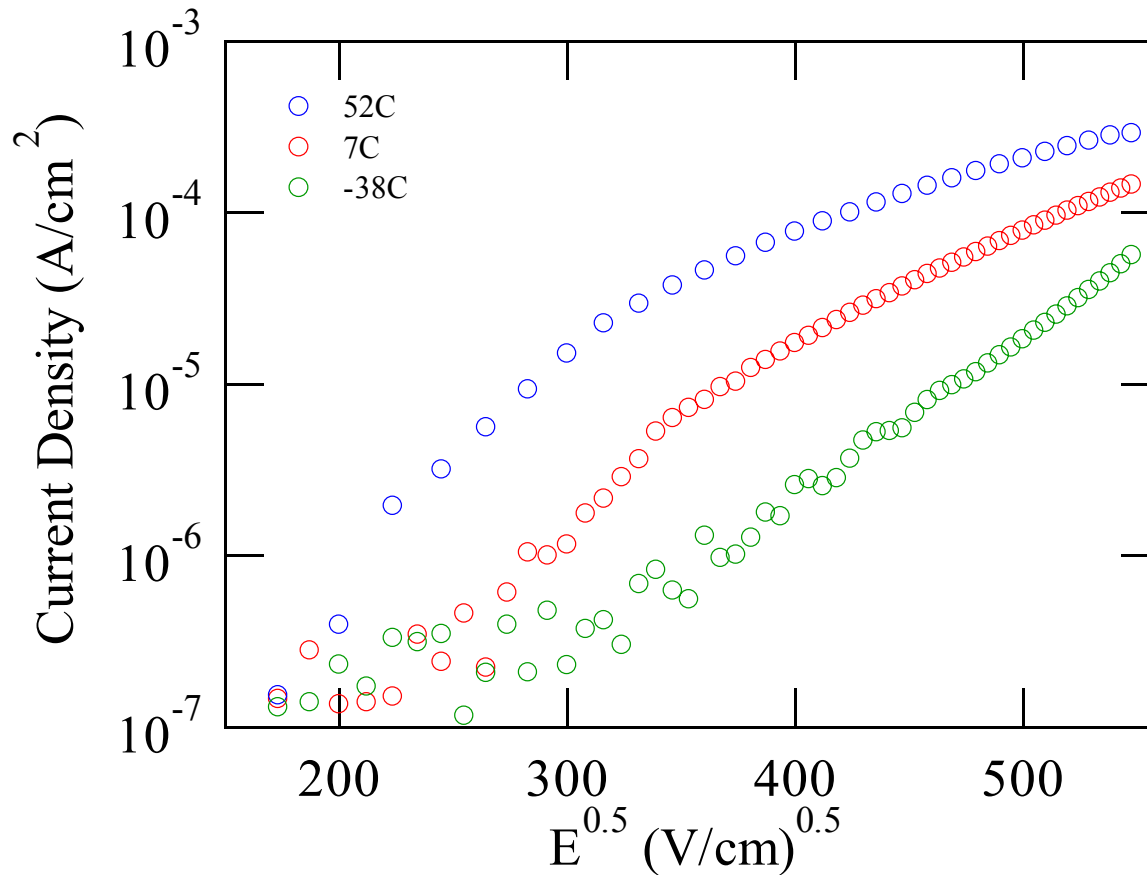
Barrier lowering



$$\Delta_{SR} = \Delta - E_{SR},$$

$$E_{SR} = \sqrt{\frac{V}{L} \frac{e^2}{4\pi\epsilon_0\epsilon}}$$

Field dependence of injection



Low mobility

RICHARDSON-SCHOTTKY EFFECT IN INSULATORS*

P. R. Emtage and J. J. O'Dwyer†

Westinghouse Research Laboratories, Pittsburgh, Pennsylvania

(Received 31 January 1966)

The Richardson-Schottky formula for thermionic emission from a metallic cathode into the conduction band of an insulator is frequently¹ stated as

$$J_S = \frac{4\pi em(kT)^2}{h^3} e^{-(\varphi_0 - \Delta\varphi)/kT}. \quad (1)$$

In this expression φ_0 is the work function, and the Schottky term is given by

$$\Delta\varphi = (e^3 F_c / \epsilon)^{1/2}, \quad (2)$$

where ϵ is the dielectric constant, and F_c the

field strength immediately in front of the cathode. It has recently been pointed out by Simmons² that this expression is invalid when the mobility of the electrons in the dielectric is low, for if one determines the density of current carriers in the insulator, n , from the relationship

$$J = ne\mu F, \quad (3)$$

one may then find that n becomes so large that back-diffusion from the dielectric to the metal will occur. Unfortunately Simmons's discus-

$$J = N \cdot e \cdot \mu \cdot E \cdot \exp(-\varphi_B/kT)$$



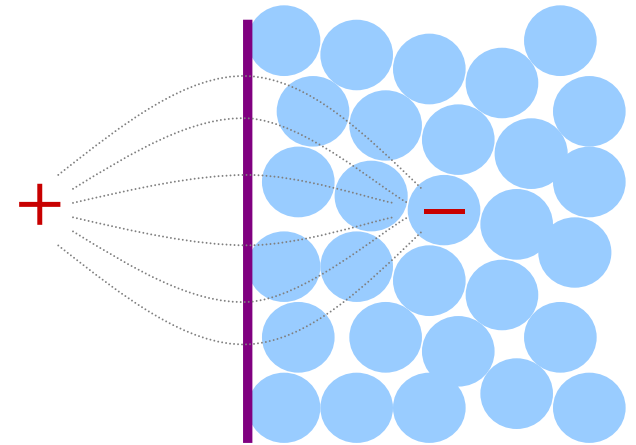
Recombination with image charge

$$J = C \exp(-\phi_B/kT) - en_0 S(E)$$

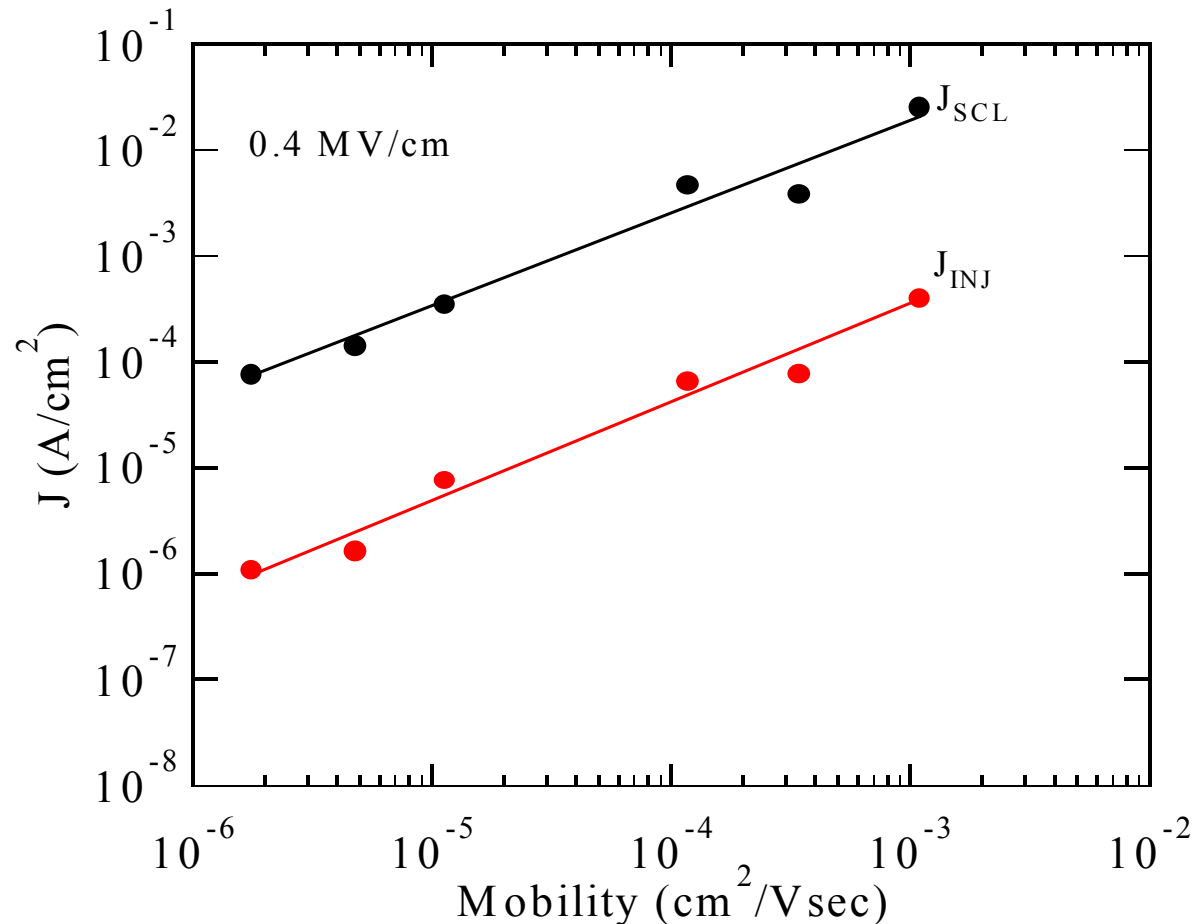
- Surface recombination as a hopping process in the image charge potential.
- No current flow at zero field.

$$C = 16\pi\epsilon\epsilon_0 N_0 (kT)^2 \mu / e^2$$

$$S(0) = 16\pi\epsilon\epsilon_0 (kT)^2 \mu / e^3$$



Dependence of injection on mobility

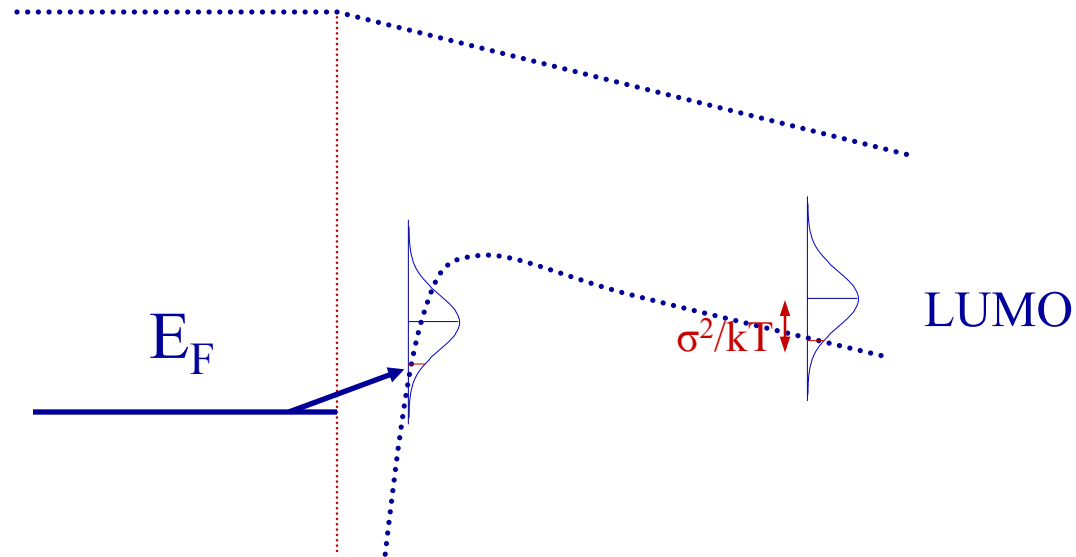


Y. Shen, M.W. Klein, D.B. Jacobs, J.C. Scott,
and G.G. Malliaras, *Phys. Rev. Lett.* **86**, 3867 (2001).



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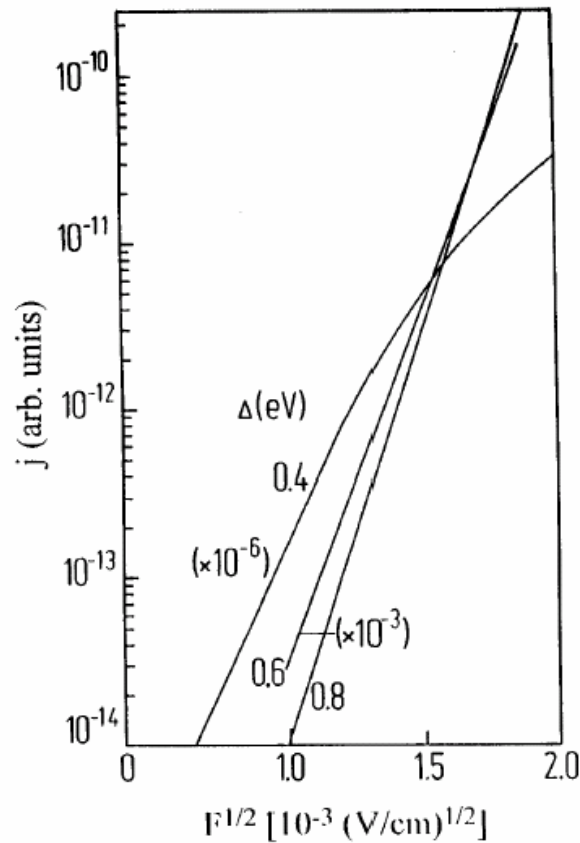
Gaussian disorder



$$J_{Inj} = e \cdot \nu \cdot \int_a^{\infty} dx_0 \left[\exp(-2 \cdot \gamma \cdot x_0) \cdot w_{esp}(x_0) \right] \cdot \int_{-\infty}^{\infty} dE [Bol(E) \cdot g(U(x_0) - E)]$$

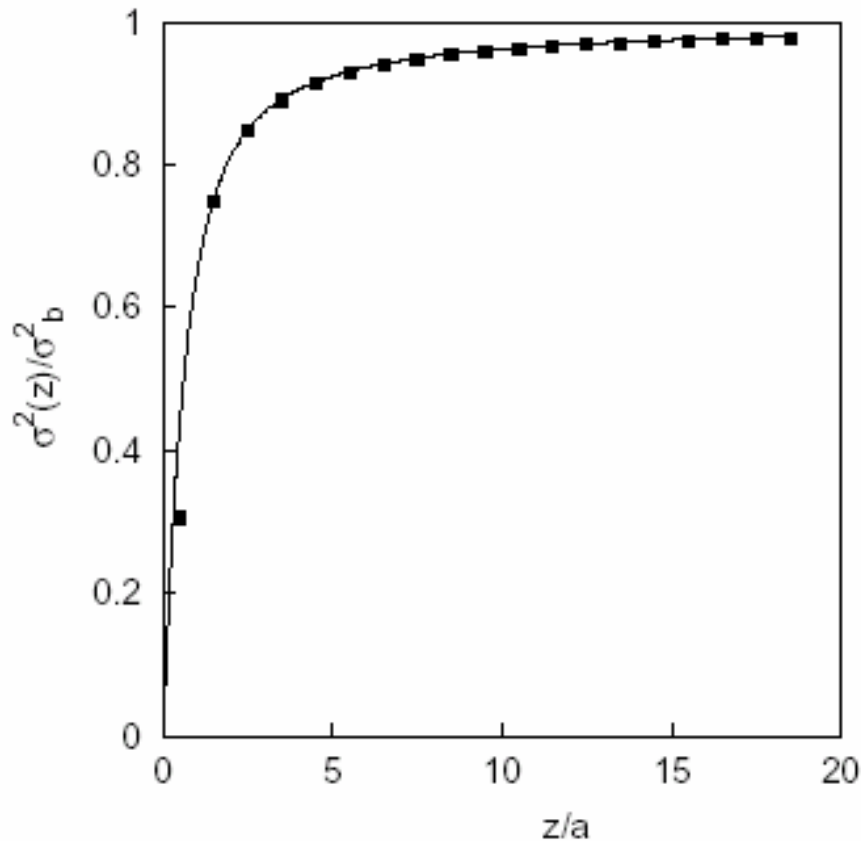
First hop is the rate limiting step

Gaussian disorder (II)



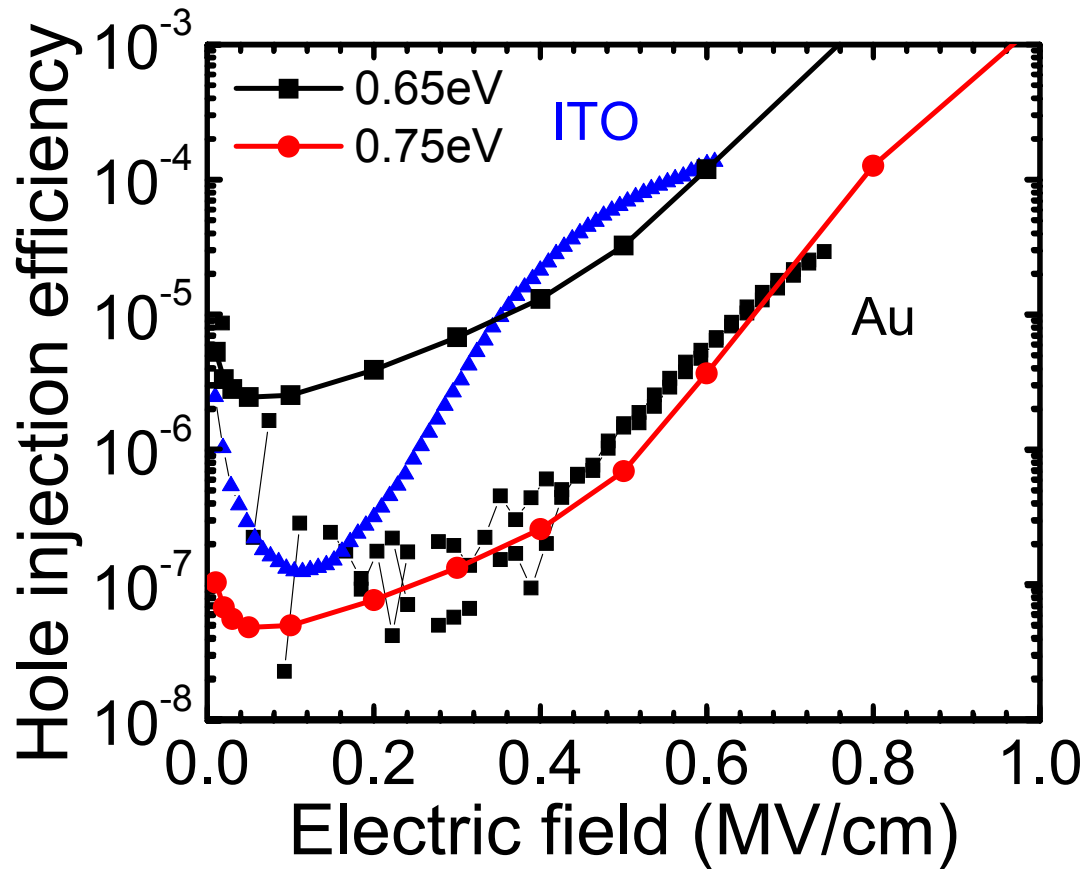
Field
dependence
resembles
thermionic
emission

Gaussian disorder (III)



Energetic
disorder due to
charge-dipole
interactions
different at the
interface

Gaussian disorder (III)



Charge injection

Mechanism:

Thermionic emission

$$J = A \cdot \exp(-\phi/kT)$$

Tunneling

$$J = A \cdot E^2 \cdot \exp(-B \cdot \phi^{3/2}/E)$$

First order corrections:

Barrier lowering

$$J \sim \exp(E^{0.5})$$

✓

Recombination with image force

$$J \sim \mu$$

✓

Disorder:

Gaussian disorder

$$J \sim \exp(E^{0.5})$$

?



Opportunities

- There is rich physics to be explored in organic light emitting diodes
- Conjugated polymers that show ideal transport characteristics and high mobilities have become available
- Charge injection in TFB is poor (and poorly understood) – opportunity for major improvements in OLED performance



Challenges

- Picture of metal/organic interfaces from spectroscopy is only now getting incorporated in injection models
- Injection expected to be spatially inhomogeneous due to correlated disorder
- $J \sim \exp(E^{0.5})$ ubiquitous, temperature range rather small
– need other tests for theories



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